

Indian Institute of Science Education & Research Pune

Magnetism in Non-Magnetic Materials

by

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*A thesis submitted in fulfillment of the requirements for
the degree of Master of Science*

under supervision of

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Scientific Research**

Dedicated to my Parents.



You've never left the existence before nor will in the future. The one who leaves is somebody else; the body leaves it's just an encapsulation, mind too leaves its microencapsulation. In between these two, there resides a master who never takes birth and never leaves the existence; he is eternal.

- Osho

Certificate

This is to certify that this dissertation entitled “**Magnetism in Non-magnetic materials**” towards the partial fulfilment of the BS-MS dual degree programme at the Indian Institute of Science Education and Research, Pune represents study/work carried out by “Kshitij Narendra Karanjekar at Jawaharlal Nehru Centre of Advanced Scientific Research” under the supervision of “Prof. A Sundaresan, Professor, Chemistry and Physics of Materials Unit” during the academic year 2018-2019.

Kshitij N. Karanjekar

Prof. A. Sundaresan

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Declaration

I hereby declare that the matter embodied in the report entitled "Magnetism in Non-Magnetic Materials" are the results of the work carried out by me at the Chemistry & Physics of Materials Unit, Jawaharlal Nehru Centre of Advanced Scientific Research, under the supervision of Prof. A Sundaresan and the same has not been submitted elsewhere for any other degree.



Kshitij N. Karanjekar



Prof. A. Sundaresan

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Abstract

The origin of magnetism in otherwise nonmagnetic materials sounds interesting. Prof. Sundaresan et al. have already shown that the ferromagnetism exists in nanoparticles of MgO, GaN, CdS, Al₂O₃, ZnO, CeO₂, In₂O₃, SnO₂, etc. due to defects on their surface. In this project, the main focus is on SiO₂ & BaTiO₃. SiO₂ is the most abundant oxide on the earth and biocompatible as well so if we can induce magnetism in it; it can be extensively used (instead of other nonbiocompatible magnetic particles) in biomedical applications while BaTiO₃ nanoparticles can also be used as vectors in nanomedicine.

In this project, we have successfully synthesized BaTiO₃ nanoparticles and different phases of SiO₂ nanoparticles. Different characterizations such as SQUID, photoluminescence and small angle neutron scattering (SANS) were essential for understanding the origin of ferromagnetism. The highest magnetic moment obtained was 0.006 emu/g in both SiO₂ & BaTiO₃. SANS data of BaTiO₃ shows that ferromagnetism indeed is coming from the surface oxygen vacancies. The reduction of magnetic moment after annealing in the air also shows that oxygen vacancies are primary source of ferromagnetism.

Chapter 1: Introduction.

1.1 Introduction.

Ferromagnetism in conventional materials arises due to the presence of unpaired electrons. However, recently it is evident that even in the absence of such unpaired electrons various oxides, nitrides, and chalcogenides (e.g.: CeO₂, Al₂O₃, NbN, CdS, etc.) show ferromagnetism at room temperature [1-3]. In the case of oxides, it has been established that only the nanoparticles behave ferromagnetically, their bulk counterpart being diamagnetic [2]. The source of the ferromagnetism has been postulated due to exchange interactions between localized electron spin moments. This might be due to the influences of the oxygen vacancies at the surface of nanoparticles [4]. However, this case should not be confused with that of dilute magnetic oxides which are wide bandgap materials that become ferromagnetic only when doped with cation having partially filled d orbital [5]. The ferromagnetism shown by thin films of HfO₂ is speculated to be a consequence of lattice defects and oxygen vacancies (induced due to preparatory conditions), and since neither Hf⁴⁺ nor O²⁻ are magnetic the ferromagnetism is termed as d⁰ ferromagnetism [6]. These types of materials can find their application in spintronics. Spintronics is the area in which along with the charge dependent electronics, spin degree of freedom can also be utilized to facilitate the manufacture of high-end devices with the increased speed of data processing and low electric power consumption [7]. The ferromagnetism observed in CaB₆ even though the constituent atoms Ca & B are nonmagnetic is surmised to be due to defects which may have formed during the synthesis of CaB₆[8]. The ferromagnetism in CaB₆ was speculated to be due to Fe based impurities which may be in the phase of Fe-B [9]. But the studies done in the single crystalline CaB₆ with a couple of different purities of Boron has ruled out the fact that the ferromagnetism is originated from Fe based impurities [10]. The role played by these surface oxygen vacancies of nanomaterials in their origin of ferromagnetism and how it differs from conventional ferromagnetic materials is to be understood in great detail.

The oxide of current interest is SiO₂. The possibility of magnetic properties in the SiO₂ may open new avenues in spintronics as well as in the field of Nano medicine. The biological applications of SiO₂ are high of interest because of its biocompatibility. It is a consensus that the properties of nanoparticles are very much different than their bulk counterparts. Similar is the case with silica. The magnetic nanoparticles which have high saturations (i.e., up to 80 emu/g) are required for

the Magnetic Resonance(MR) contrast agents for MR diagnosis of cancer [11]. The Fe_3O_4 nanocrystals have already been observed as good intracellular labeling agents and good MR trackers [12].

The magnetic materials with high magnetic saturations and their oxides are considered mainly for therapeutic applications [11]. In biomedical applications the Nano phase is always preferred, because of smaller size it is easier for nanomaterials to travel in veins and arteries. But along with having benefits of Nano level dimensions, there is a probability of these particles getting penetrated through membranes only to pass through the bloodstream to other organs [13]. Fe_3O_4 nanoparticles are found to induce an inflammation, cellular poisoning and respiratory damage [13]. So the magnetic materials in such applications are required to be biocompatible. SiO_2 nanoparticles can be better candidates for such biomedical applications due to their biocompatibility. Experimental studies on the effects of SiO_2 upon cellular uptake doesn't show any cytotoxic activity up to 2.5nM of concentration [14]. While the mesoporous silica (silica containing the high specific surface area and pores) are effective in drug delivery [15], So these SiO_2 nanoparticles if made somehow magnetic then they can replace the otherwise toxic materials in biomedical applications such as hyperthermia, MR labeling agents, magnetic separation of biological molecules, cell labeling, drug carriers, radionuclide and gene delivery as well as tissue engineering and magnetic bio sensing [16].

But as is known that the commercial pyrogenic SiO_2 Nano powder (NP) is diamagnetic [17]. But the synthesis of various phases of SiO_2 can be carried out to explore the possibilities of magnetic properties. Nano powders produced by Vladislav et al. shows d0 magnetism [17]. The studies which were done by Gregory Kopnov et al. have shown that the etching of the silicon wafers by HF and KOH have shown relatively large magnetic moment. Their studies suggest that the magnetism is co-related to surface roughness; the exact correlation is unknown. The precise mechanism of this magnetism is currently an open question [18]. But there is another phase of SiO_2 which is that of SiO_2 Nano spheres which are also shown to exhibit weak ferromagnetism. The hole-bridge structures of which were obtained by etching the already sintered Nano spherical array with HF. In the work of Xing Wang et al. the origin of weak ferromagnetism in this phase was assumed to be none other than oxygen vacancies. They have speculated that HF-etching may have contributed to weak ferromagnetism in these Nano spherical arrays [19].

Another oxide of interest is BaTiO_3 . The nanoparticles of which show ferromagnetic order and also hold a promising future in the biomedical field [20-21].

In this project, the synthesis of BaTiO₃ nanoparticles & various phases of SiO₂ nanoparticles will be carried out. The characterization of which will be done by XRD. The size and shape will be studied using scanning electron microscopy(SEM). The magnetic moment will be measured with SQUID.

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Chapter 2: Sol-Gel Synthesis.

2.1 Background.

Since in this thesis we are mainly concerned about nanoparticles we will discuss general methods of synthesis of nanoparticles. There are two main approaches which are top-down & bottom-up approaches [1-2] former being costlier than the latter. In the bottom-up approach the atoms or molecules are allowed to self-assemble naturally or through an externally driven force in a chemical reaction to form more organized Nano structures while in top-down approach starting with a larger piece finer and finer structures are made by using various tools and methods [2]. Here we are mainly concerned about the bottom-up approach that too particularly sol-gel and precipitation methods.

2.2 Bottom-up approach.

The bottom-up approach is again divided into several methods such as molecular self-assembly, molecular beam epitaxy, sol-gel, etc. [3-13] In molecular self-assembly molecules are allowed to form definite structures naturally, i.e., without external driving force [14]. Molecular self-assembly is much more relevant to biological entities such as peptides & proteins and some extent to polymer science and materials engineering [15-16]. Molecular beam epitaxy is a layer by layer deposition of single crystals to form the thin film under very high vacuum [17]. It is used for the synthesis of single crystal thin films, quantum wells & superlattices [18]. Sol-gel synthesis is the novel method for synthesis of metal oxide nanoparticles or nanocomposites. It is novel as in it has control over the surface properties of as-synthesized nanoparticles [19]. It also is the favored method for studying surface modification by coating a layer, surface doping, coating with organic molecules, etc. [20-24].

2.3 Sol-gel method.

A sol-gel method is a method which uses a solution containing metal compounds such as metal alkoxides as a precursor for the oxide to be prepared, alcohol as a solvent, H₂O as a hydrolyzing agent, any acid or base as a catalyst [25]. A sol is generally referred to as solvent hosting colloidal particles. The process starts with a dissolution of metal precursors which is more often

a 'metal-alkoxide' in a solvent to form a sol [26]. Except for metal alkoxides such as silicon alkoxide, hydrolysis of other metal precursors yields particulate materials rather than undergoing polymerization; hence much of the sol-gel works are based on silica glass & silica gels [27]. After formation of sol, gelling (formation of gel) takes place by polymerization which generally starts from the surface and proceeds in bulk [28]. So in the following three subsections, we'll discuss three aspects of our reactions, i.e., hydrolysis, condensation, gelation & drying.

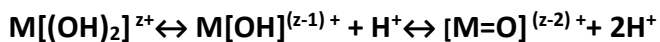
2.3.1 Hydrolysis.

There are two types of precursors generally used one is an aqueous solution of an inorganic salt, and other is a metal organic compound [29]. The most commonly used precursors are metal alkoxides given their high reactivity towards nucleophilic species such as H₂O [30]. When the metal cation is introduced as a salt in the water it is solvated by water molecules as shown below in fig 2.1; the electrons are transferred from the bonding orbitals of H₂O to the transition metal cation which creates a partial charge on H atom making the molecule acidic [29].



fig 2.1 Solvation of transition metal cation.

So, the following equilibrium exists which is dependent on charge transfer and in turn acidity of the H₂O molecule [29].



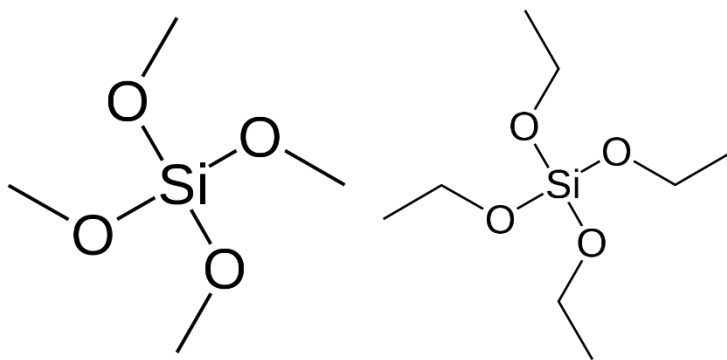


fig 2.2 The chemical structures of TMOS & TEOS respectively.

Precursors typically used for the synthesis of silica nanoparticles are tetramethyl orthosilicate (TMOS) and tetraethyl orthosilicate (TEOS) though we'll use TEOS as our main precursor. Both the reagents come under the category of silanes with four alkoxy groups attached to the Si atom as shown in fig 2.2. These precursors are generally soluble in organic solvents such as ethanol & methanol etc. here we've mostly stuck to the ethanol. The hydrolysis of silanes requires catalyst which is generally of two types acidic or basic, the reactions under both catalysts are shown in fig 2.3 & 2.4 respectively [30].

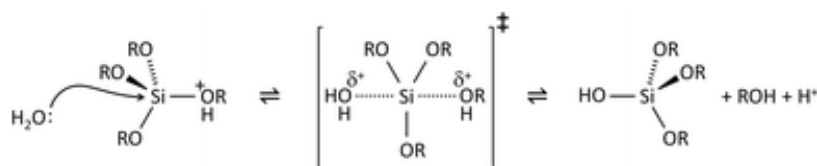


fig 2.3 Acid-catalyzed Hydrolysis.

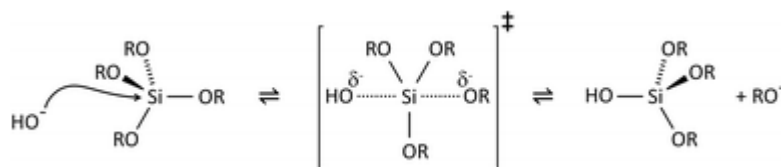


fig 2.4 Base-catalyzed hydrolysis.

2.3.2 Condensation.

Condensation can be proceeded by two modes either nucleophilic substitution (S_N) or nucleophilic addition (A_N) as following [25];

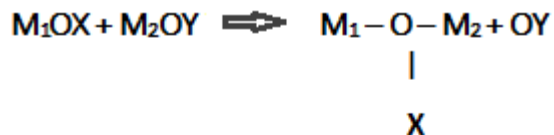


fig 2.5 Condensation by S_N & A_N modes.

2.3.2.1 Olation.

One of the ways by which condensation takes place is by the formation of hydroxyl bridge between two central metal atoms known as olation [32]. One of the ways the olation can take

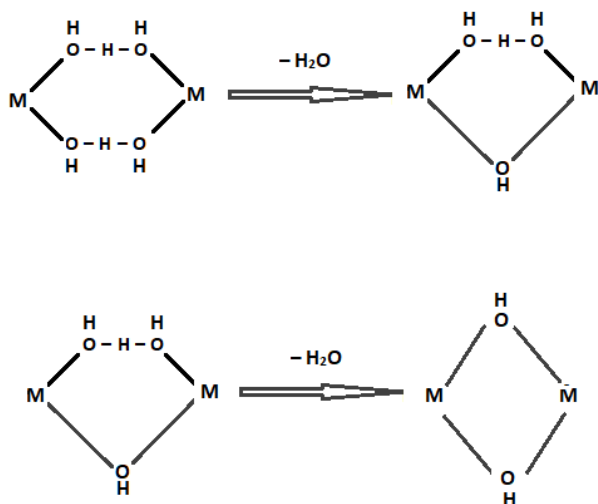


fig 2.6 Olation by H_3O_2 bridge formation [25].

place is by the formation of H_3O_2 bridge (fig 2.6) [33-34].

But this is not the only way by which the oleation occurs there are other ways too as shown in fig 2.7 [29,25,35] The rate of ololation depends upon size, electronegativity & electronic configuration of metal cation M [36-38,52].

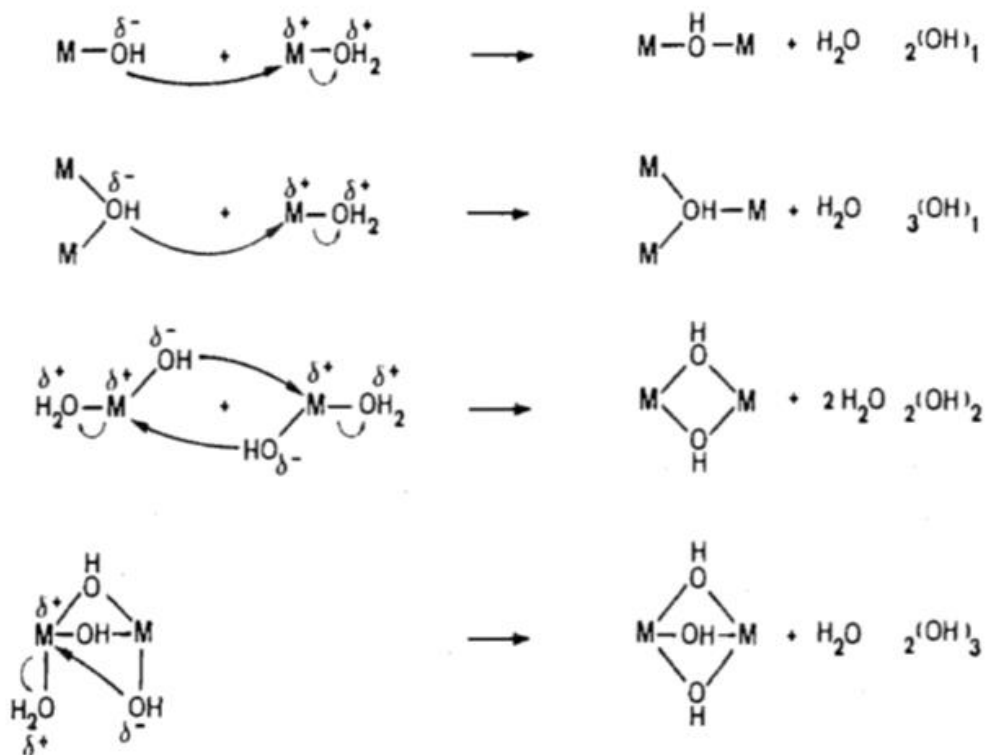


fig 2.7 Different mechanisms of Oleation [25].

2.3.2.2 Oxolation.

In oxolation, the bridge between two central metallic atoms is the -O- (oxo-) group. This type of reaction also takes place in two ways, i.e., S_N & A_N. Former is the case for coordinately saturated metals & latter is for coordinately unsaturated metals [39-41].

2.3.2.3 Condensation of silicates.

After hydrolysis the hydrolyzed silica interact with each other to form $\equiv Si - O - Si \equiv$ bonds resulting in linkage [42] as shown below (fig 2.8);

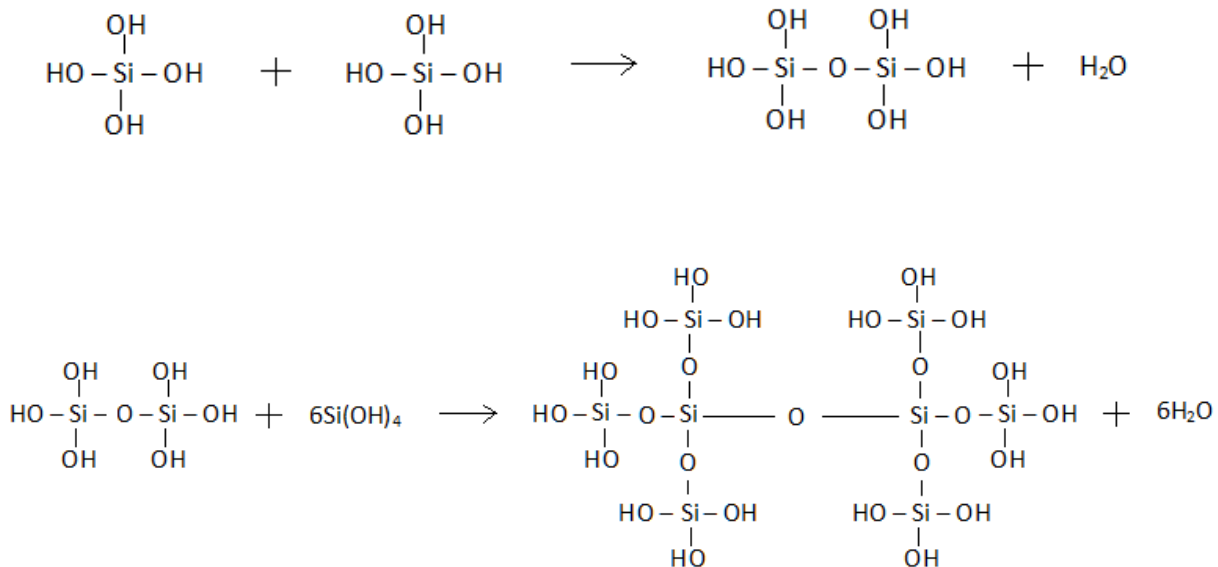


fig 2.8 Condensation.

Former reaction depicts the simplest condensation that can take place between two molecules while the latter shows the poly condensation. But it is not always the case since the silicates are most of the times never fully hydrolyzed (depending on catalyst), i.e., some –OR groups are still present so for such cases condensation takes place by producing alcohols (fig 2.9) [43].

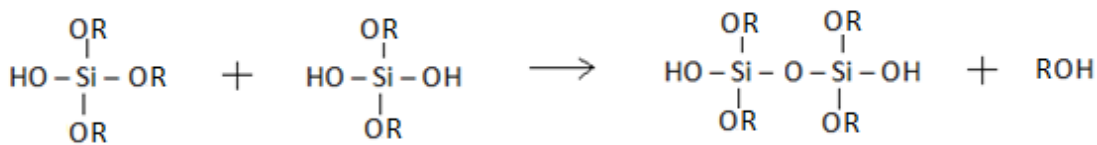


fig 2.9 Condensation.

2.3.3 Gelation & drying.

Gelation can be defined as a formation of a three-dimensional solid network (mostly polymer) in a fluid phase [31]. The turbidity of the solution increases gradually [44]. The simulation studies by using Brownian & Stokesian dynamics have shown that gelation time (t_g) depends on salt concentration & volume fraction of particles [45]. The former affects the gelation time but not

the strength of the gel network [45]. Gelation time (t_g) is the time after which the semisolid leftover can bear stress elastically [42]. In the final process of drying the semisolid matter turns into solid also called as Xerogel [44]. After drying; the solid may or may not be calcined to obtain the final product.

2.4 SiO₂.

SiO₂ is an oxide of silicon also known as silica. It is the most abundant oxide on the earth. Silica constitutes 10% of the mass of the earth's crust in the form of crystalline quartz [46]. Silica is most abundant component of sand [47]. In the biological world, it is commonly found in shells of aquatic creatures, in plants and rice husk. It has a wide variety of uses. Almost 95% of commercial production of SiO₂ ends up in construction applications [46]. It is the main constituent of glasses. Silicones are polymers of $\equiv \text{Si} - \text{O} - \text{Si} \equiv$ (siloxane) groups. Silicones are used as coatings & insulators in electrical appliances due to their low thermal & electrical conductivity [46]. It is also used as a defoamer, oil, grease, rubber, dry cleaning solvent, firestop, in medicine particularly ophthalmology, mold making, etc.

Silica is mostly found in tetrahedral co-ordination with oxygen. Silica can be either crystalline or amorphous. Crystalline silica has 14 types of polymorphs. We'll mostly be concerned of amorphous SiO₂.

2.5 Synthesis of Silica Nanoparticles.

Silica nanoparticles were first prepared in 1956 by Kolbe by mixing TEOS in a solvent containing alcohol & water but they were agglomerated & the results were only observed for very pure reagents [48]. But after systematic study of the reaction parameters involved Stober et al. successfully synthesized silica nanoparticles even without using very pure reagents and thus inspired the silica sol-gel syntheses yet to come [49]. The work of Stober was mostly extended by Bogush et al. The Nano spherical SiO₂ particle size depends on multiple factors. It depends on the concentration of TEOS, NH₄OH, H₂O, reaction temperature as well as the ratio of ethanol and aqueous ammonia [50]. The studies have also shown that the average particle size goes through a maximum as a function of water concentrations as shown in figs 2.10 & 2.11 [51]. The maxima of these curves are a function of ammonia concentration.

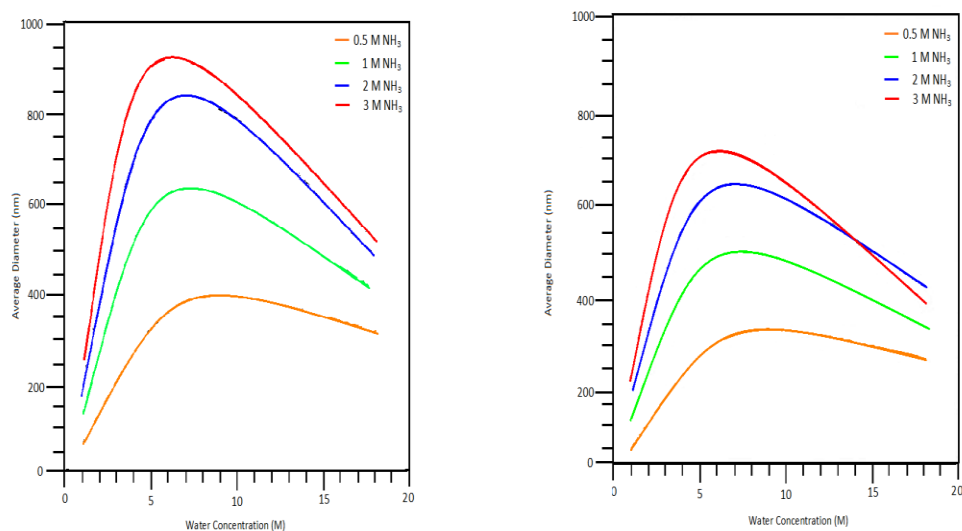


fig 2.10 Average diameter of as-synthesized SiO₂ nanoparticles as a function of water concentration at fixed 0.5,1,2,3 M concentrations of NH₃ & at fixed concentrations of 0.17 M & 0.3 M of TEOS respectively at temperature 298K.

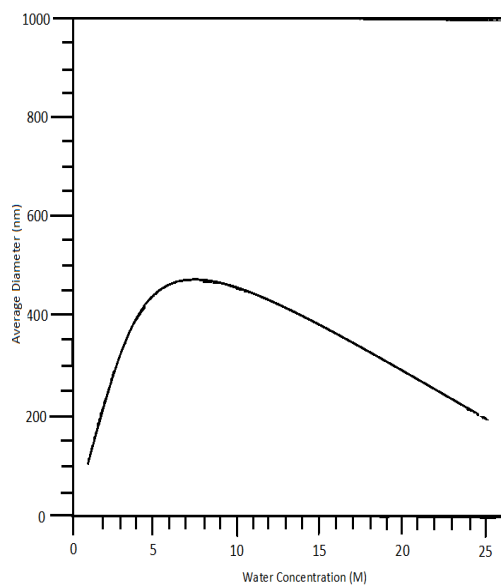


fig 2.11 Average diameter of SiO₂ nanoparticles as a function of water concentration at fixed concentrations, i.e., 0.17M & 1 M of TEOS & NH₃ respectively at 298K.

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Chapter 3: Experimental Techniques.

3.1 Overview.

This chapter is dedicated to the understanding of different characterization techniques & working principle behind them. We'll mainly discuss x-ray diffraction (XRD), scanning electron microscopy (SEM), fourier transform infrared spectroscopy (FT-IR), photoluminescence spectroscopy (PL), and ultraviolet spectroscopy (UV).

3.2 X-ray diffraction (XRD).

XRD is a characterization technique which is extensively used in solid-state sciences. X-ray diffraction in crystals was first observed by Max Van Laue et al. [1] & the development of bragg equation as well as usage of this phenomenon to characterize the Zinc Blend structure was done by both W. H. Bragg & W.L. Bragg a father & son duo for which they were awarded Nobel prize for physics in 1915 [2]. On the electromagnetic spectrum, the range of wavelength from 0.01 nm to 10 nm is termed as x-ray. X-rays have a relatively high penetrating power as is evident for their application in medical imaging of bones [3]. When a charged particle decelerates it emits radiation, this radiation is often termed as bremsstrahlung. So in an XRD instrument, the high voltage is applied between cathode and anode, so the electrons upon being incident on anode decelerate and emit x-rays, above certain threshold voltage molybdenum will show two characteristic lines $K\alpha$ & $K\beta$ which are evident as two sharp peaks under the continuous background of bremsstrahlung [4-5]. The relation between the x-rays and the spacing between lattice planes is given by Bragg's law i.e.

$$n\lambda = 2d \sin\theta$$

Where n is an integer, λ is a wavelength, d is interplanar distance, θ is the angle of diffraction. The more rigorous mathematical treatment by Max von Laue using reciprocal lattice is also equivalent to the Bragg's law but will not be discussed here. There are two main ways to produce electrons, one way is by ionization of the gas, and another way is by deceleration of electrons on the target by applying voltage [6]. As can be seen in fig 3.2 the target acts as an anode while

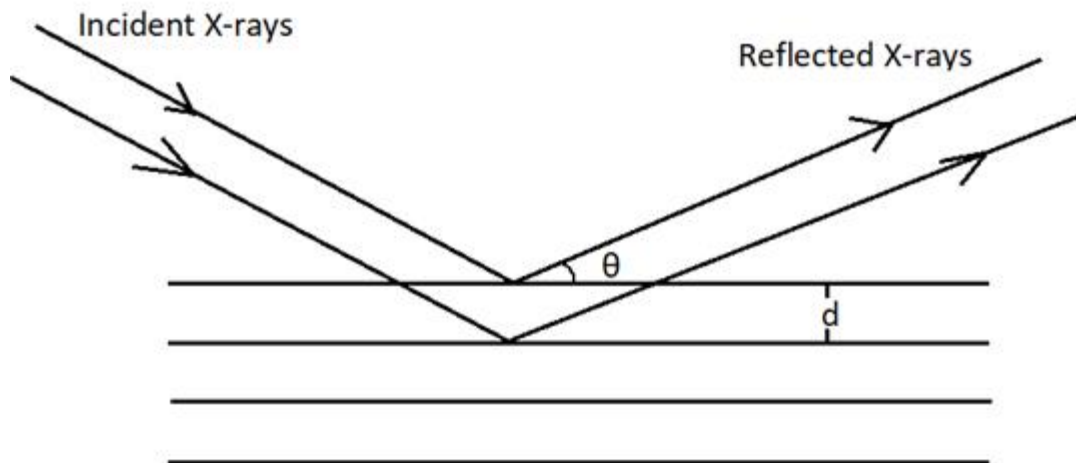


fig 3.1 Incident X-rays are diffracted from lattice.

there is a tungsten filament which acts as a cathode & is pumped with a current of 3 ampere to provide the heating for ejection of electrons while the target receives the electron [6]. The target gets heated so to cool it down water is flowed in an adjacent channel [6]. The x rays are scattered in all the directions and are used for the diffraction.

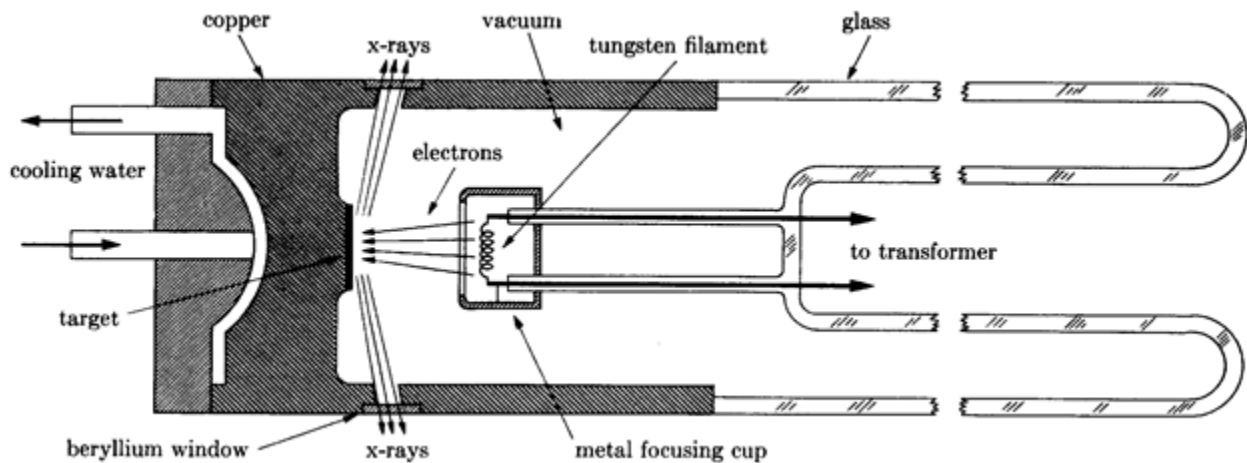


fig 3.2 Cross section of X-ray tube.

3.3 Scanning electron microscopy.

Scanning electron microscope (SEM) uses electrons to probe the microscopic structure of matter, unlike optical microscope which uses light radiation. The resolving power of SEM is higher

because of the smaller wavelength associated with electrons (0.005 nm) rather than light (200-250nm) [7]. The electrons upon incidence on sample interact with it through various interactions which involve emission of photons as well as the scattering of electrons [8]. Among these, we will mainly discuss back scattered electrons and secondary emission of electrons. Secondary electron emission is the emission of the electron when radiation is incident on a metal surface which was first observed by Lenard et al. [9]. The secondary electrons (SE) are mainly the valence electrons that are inelastically scattered upon the impact of primary beam of electrons [10]. The secondary emission electrons are collected by a detector, and the topography of material is scanned and imaged on CRT display (nowadays LCD) [11]. The SE have energies which are usually less than 50 eV with mean free path approximately 1 nm [12]. The electrons which are emitted are not only SE but also the inelastically back scattered & elastically scattered, i.e., reflected electrons [12]. An Everhart-Thornley detector is used for detecting the secondary as well as back scattered electrons [13]. Everhart-Thornely detector is a cylindrical brass whose end is fitted with copper mesh with positive voltage so upon incidence the electrons are accelerated to this copper mesh only to find themselves incident on scintillator which gives rise to number of photons and these photons are incident on the photomultiplier tube which thanks to photoelectric effect gives rise to multiple electrons which are then raster scanned on cathode ray tube (CRT) [14]. But the instrument which we have used has an LCD. Fig 3.3 depicts the construction of the scanning electron microscope [15]. The electron gun emits the electrons which are to be accelerated by so-called electromagnetic lenses for focusing [16]. These electrons are incident on the specimen which is kept in high vacuum [17]. And the secondary electron detector collects the SE [18].

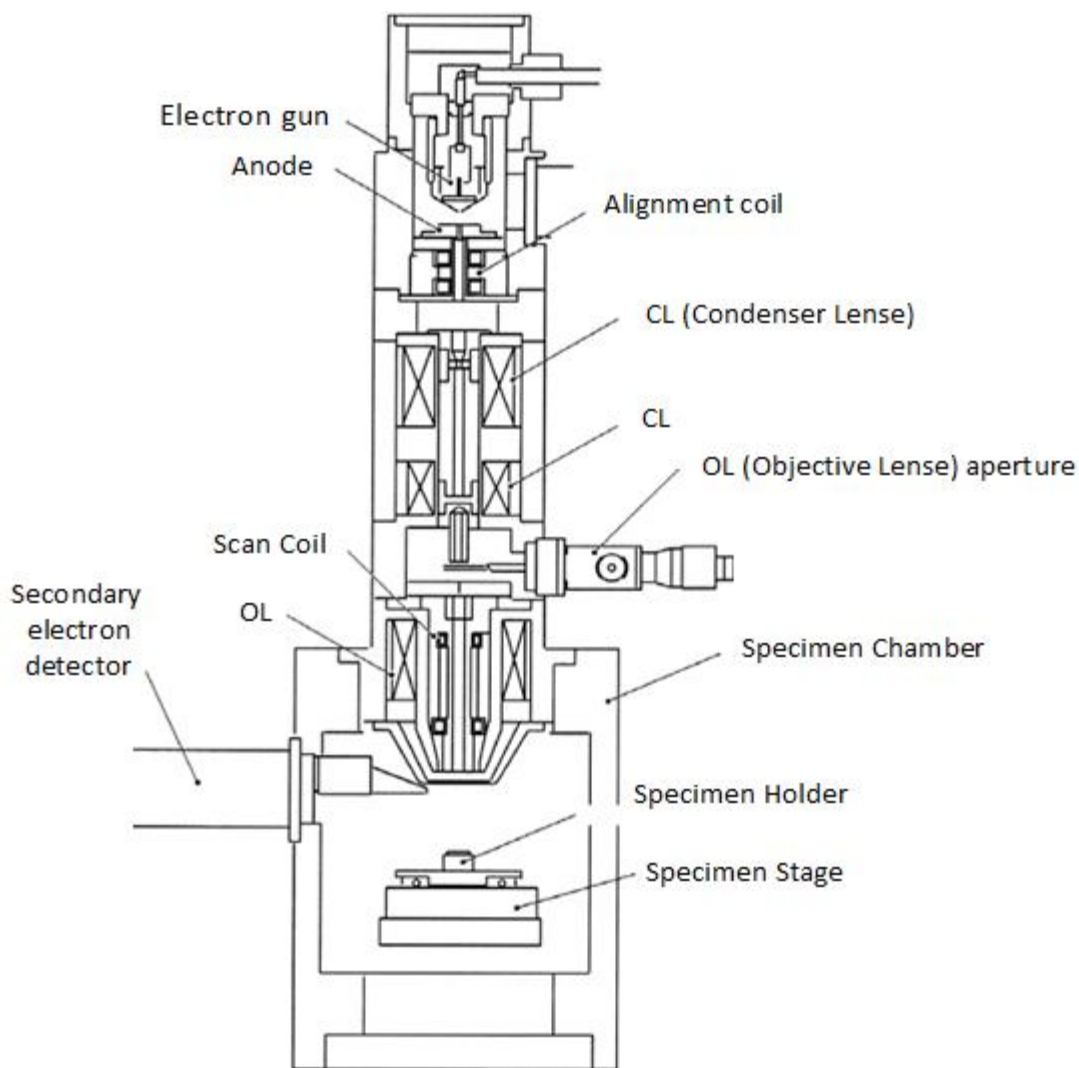


fig 3.3 Anatomy of Scanning Electron Microscope.

3.4 UV visible Spectroscopy.

UV visible spectroscopy operates in UV & visible range (200-800 nm) [19]. The parameters which are measured are absorbance in case of liquids and reflectance in case of solids. In the case of solids, this spectroscopy is referred to as diffuse reflectance spectroscopy. Bohr's model defines atoms as consisting of n number of shells (each having its specific energy) in which the electrons revolve around the nucleus, among these, there is also a state which is called as the ground state from which the electrons go to excited state upon incident radiation [20]. In the case of molecules σ , π and non-bonding electrons are present [21]. Upon incident radiation transition

from filled orbital to empty orbital (usually antibonding orbital) can occur [21]. These transitions require high energy to be excited, i.e., lower wavelengths (near UV region) [21]. When light passes through a solution, its initial intensity (I_0) decreases to some final intensity (I) because of absorbance which according to Beer Lambert's law is directly proportional to concentration & path length of light through solution [22].

$$A = \alpha l c$$

Where A is absorbance, l is path length of light through solution & c is the concentration of solute. Transmittance (T) is given in terms of I & I_0 as [23];

$$T = \frac{I}{I_0}$$

& related to absorbance (A) as [23];

$$A = \log_{10}\left(\frac{1}{T}\right)$$

Since our samples are powder, we will be mostly concerned about reflectance. To measure reflectance, we will utilize the Diffuse reflectance UV visible spectroscopy technique. Specular Fresnel reflection is the one in which angle of incidence is the same as the angle of reflection. When light is incident on the irregular sample surface, the individual Fresnel reflections give rise to an apparent phenomenon (where the angle of reflection doesn't depend on the angle of incidence) called diffuse Fresnel reflection as shown in fig 3.4 [24-25]. The widely accepted theory explaining Diffuse reflectance spectra is Kubleka-Munk (KM) theory. A solid sample consists of many particles/grains randomly distributed when light radiation is incident on the sample apart from diffuse Fresnel reflectance it can also undergo refraction, and after refraction, it may undergo many reflections or diffractions and exit the sample surface [26]. This type of radiation is termed as Kubleka-Munk reflectance as shown in fig 3.5 [26].

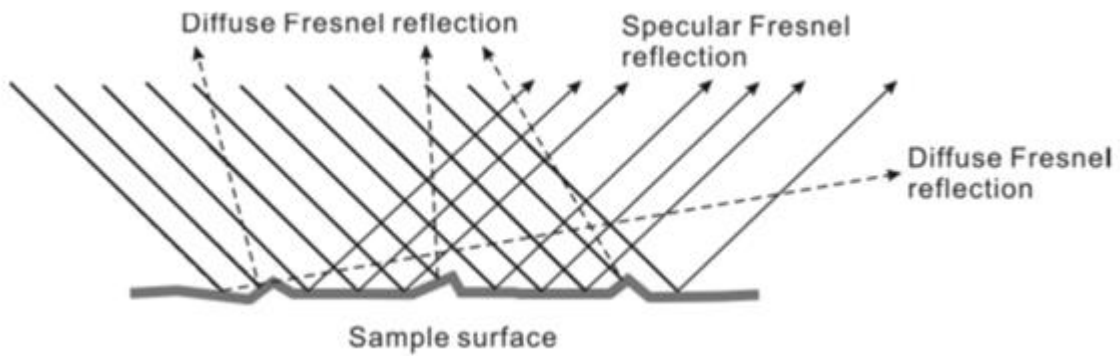


fig 3.4 Difference between Diffuse Fresnel & Specular Fresnel reflections.

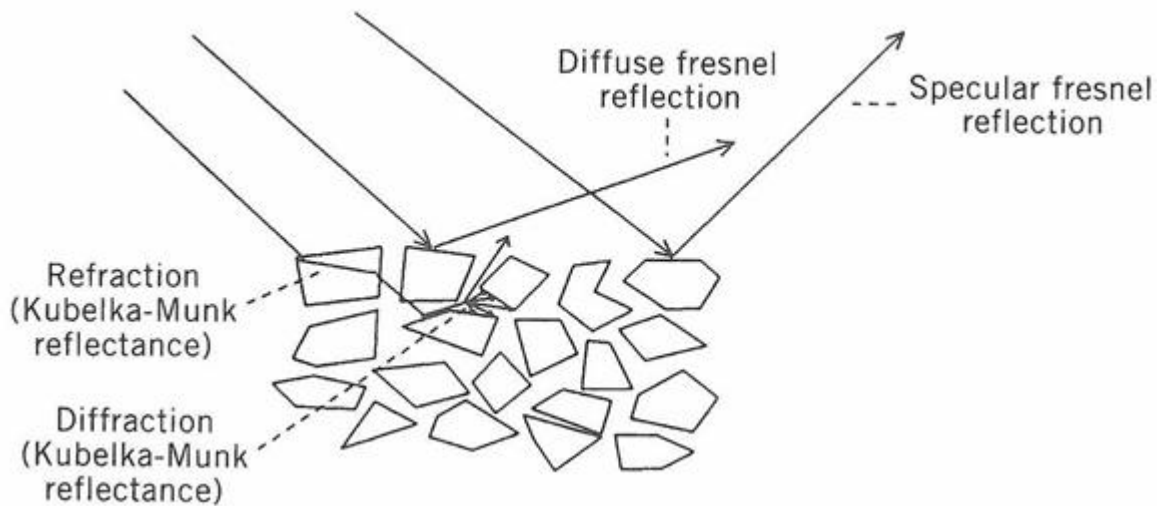


fig 3.5 KM reflectance in contrast to Specular & Diffuse Fresnel reflections.

KM theory is mainly a two flux theory, i.e., it deals with the radiation incident from two opposite directions on the layer of a material [27]. There are two components to this theory which are absorption coefficient (K) & scattering coefficient (S) as is apparent from KM function given as [28];

$$\frac{K}{S} = \frac{(1 - R_{\infty})^2}{2 R_{\infty}}$$

Where R_{∞} is absolute remittance. The ∞ symbol denotes that the sample/material thickness should be infinite, i.e., none of the light radiation should completely penetrate the sample. The assumptions of KM theory are as follows [26,29];

1. Fresnel reflection is to be ignored.
2. The constituent grains/particles are smaller than the thickness of the sample.
3. The light radiation does not penetrate the sample.
4. The sample area should be much larger than the thickness.

Now we will discuss the instrumentation of the diffuse reflectance spectrophotometer. An integrating sphere is used for data acquisition. An integrating sphere is a hollow spherical container coated with a nearly 100% reflecting surface from inside to diffuse the input light radiation uniformly on the inside surface [30]. The reference sample (one with nearly 100% reflectance) is first situated in a holder as shown in fig 3.6 for the auto-correction then the sample of interest is placed in the holder to determine the relative reflectance [26].

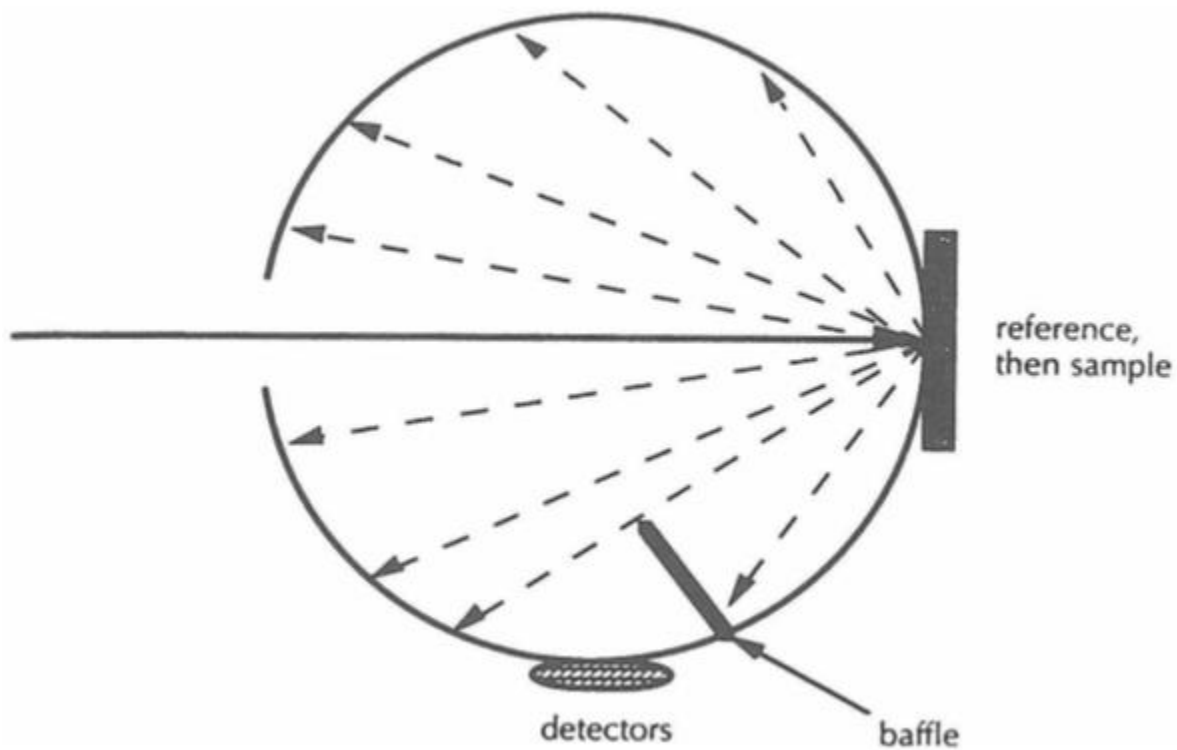


fig 3.6 Integrating sphere; an optical component to diffuse the incoming light radiation.

The detector detects the radiation intensity & we thus obtain the reflectance spectra of the sample.

3.5 Fourier Transform Infra-Red Spectroscopy.

The energy of the molecule consists of electronic, rotational, vibrational & translational energy [31]. The electronic transitions occur at UV-visible region, rotational transitions at the microwave region while the vibrational at infra-red region of the electromagnetic spectra [31]. The infra-red region occurs between 700nm – 1mm [32]. So study in this region is important to understand the types of bonds present in a given compound. The vibrational and rotational motions of the molecules are the ones responsible for the infra-red spectra [33]. In a classical mechanics' point of view, a dipole can be visualized as a simple harmonic oscillator in which the force between the two atoms follow the linear relationship with a distance (Hooke's Law) where the vibrational frequency ν is given as [34];

$$\nu = \frac{1}{2\pi c} \sqrt{\frac{k}{\mu}}$$

Where c is the speed of light, k is the force constant & μ is the reduced mass of the dipole. In quantum mechanical domain when the Schrodinger equation is solved using simple harmonic oscillator the obtained energy eigen values are quantized [35];

$$E_n = (n + \frac{1}{2})\hbar\omega$$

Where \hbar is Planck's constant, ω is the frequency of oscillation and $n = 0, 1, 2, 3, \dots$. Thus the energy difference between the two states is $\hbar\omega$ as shown in fig 3.7. But in reality, the simple harmonic oscillation is just an approximation, so anharmonicity should also be considered. The solid curve shown in fig 3.8 [37] is the one which accounts for this anharmonicity (Morse curve) as opposed to the dashed curve (simple harmonic approximation). The potential is given as [36]

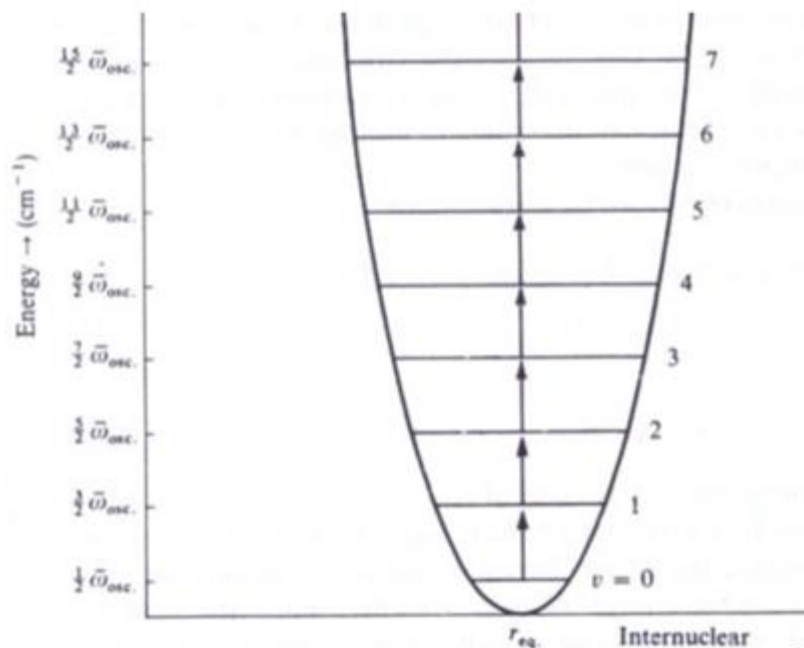


fig 3.7 The allowed vibrational transitions are shown for a diatomic molecule in simple harmonic approximation.

$$V(r) = D \{1 - \exp[-b(r - r_e)]\}^2 - D$$

Where D is the dissociation energy, b is constant, r is inter nuclear distance & r_e is the equilibrium internuclear distance [36-37]. The energy levels now are given as [37]

$$E_v = \left(n + \frac{1}{2}\right) \tilde{\omega}_e - \left(n + \frac{1}{2}\right)^2 \tilde{\omega}_e \chi_e \text{ cm}^{-1}$$

The energy levels are now contained within the Morse potential where $\tilde{\omega}_e$ is oscillation frequency & χ_e is anharmonicity constant.

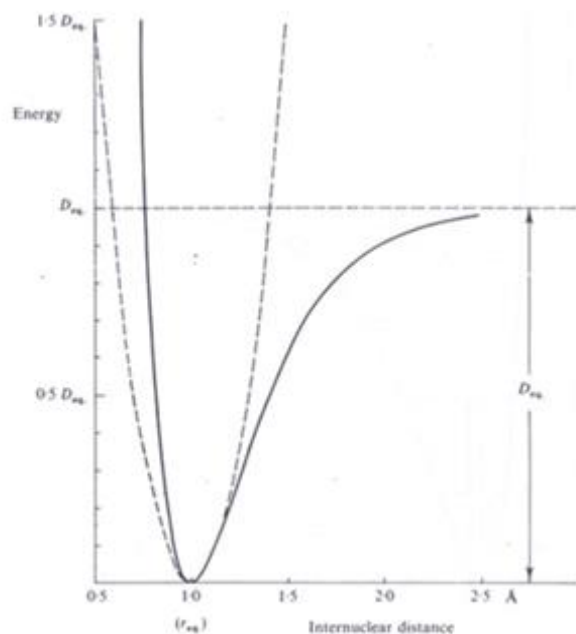


fig 3.8 The curve with dashed line is for ideal simple harmonic oscillator while the solid curve is due to Morse potential.

3.5.1 Working of FT-IR.

FT-IR is based on Michelson interferometry [38-39]. The set up consists of beam splitter, fixed & moving mirror, detector & IR source [39]. The IR source emits radiation in the direction of a beam splitter which reflects the incident beam towards fixed & moving mirrors (the splitting is 50-50%) [39]. The mirrors again reflect the radiation towards beam splitter which in turn again reflects half the beam towards detector and half towards the IR source [39]. The format in which the signal is recorded is known as interferogram which is then converted to spectrum by Fast fourier transform algorithm [39]. The stretching, bending modes are thus probed by their characteristic frequencies [40-41].

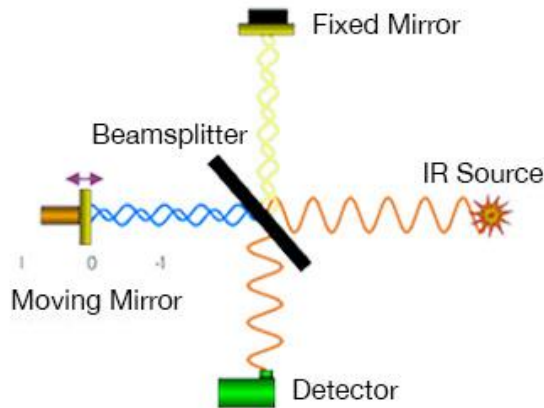


fig 3.9 Typical set up of FT-IR spectrometer [38].

3.6 Energy dispersive x-ray spectroscopy (EDS).

EDS is a technique which is used along with SEM for characterization of elements present (along with their respective proportions) in the sample [42]. There are electron distributions around the atomic nucleus which are referred to as shells, i.e., K, L, M, etc. [43]. When an electron beam is incident on a sample two types of radiations are realized: one is characteristic x-ray & other is continuum x-ray [42]. Both are the result of inelastic interactions [42]. Characteristic x rays are the ones which are emitted by atoms due to electronic transitions upon bombardment of incident electrons [44]. While the continuum x-ray/bremsstrahlung is the radiation emitted by electrons upon deceleration [45], they represent the background [42]. So characteristic x rays are the ones who give useful information about the elemental composition of the sample as can be inferred from following equation (Moseley's equation) which relates atomic number Z to frequency ν [6].

$$\sqrt{\nu} = C (Z - \sigma)$$

Where C & σ are constants. The characteristic x rays which are emitted by the sample are taken up by the detector which converts the signal into the voltage pulse which in turn is converted in digital signal to give the spectrum [46]. The data can also be represented in an image format also known as elemental mapping.

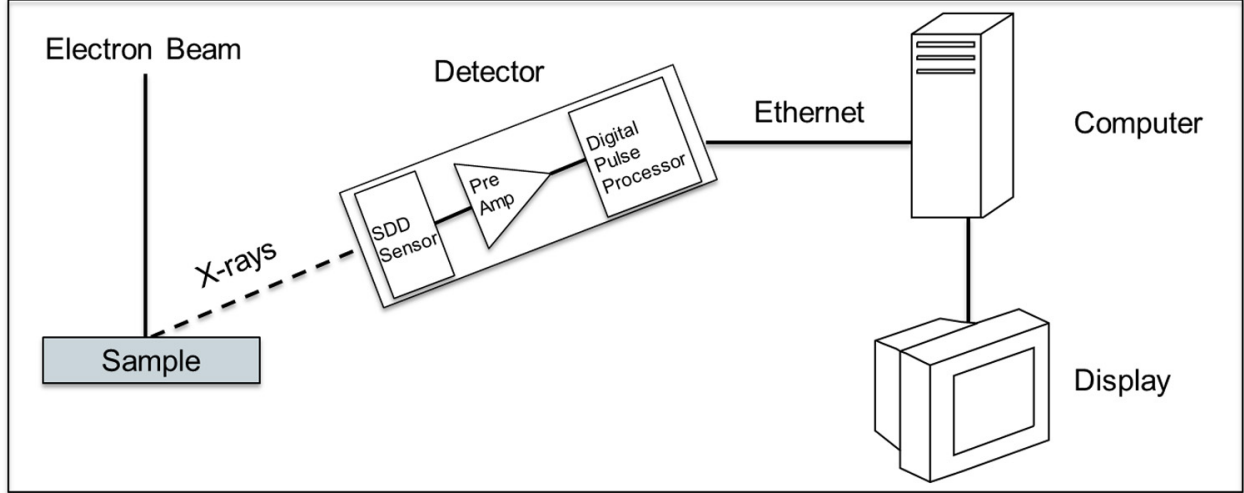


fig 3.10 Construction of EDS [46].

3.7 Small angle neutron scattering (SANS).

3.7.1 Scattering (background).

When particles are incident on some target; due to interactions with the target's potential their path is deviated. This phenomenon is known as scattering. Scattering is a useful tool in physics to probe the microscopic matter. Scattering cross section is the quantity which is determined in any scattering experiment. Scattering cross section is the number of particles scattered per unit time per unit flux of incident radiation [47]. When we are interested in detecting the number of particles in a small element of solid angle ($d\Omega$), then this quantity is referred to as differential cross section [48]. In the case of classical Rutherford scattering the differential cross section is given by [49];

$$\frac{d\sigma}{d\Omega} = \frac{Z^2 e^4}{(4\pi\epsilon_0)^2 m^2 v_0^4 \sin^4\left(\frac{\theta}{2}\right)}$$

Where Z is a nuclear charge of the target atoms, e is the electronic charge, ϵ_0 is the permittivity of free space, m is the mass of incident particles, v_0 is the velocity of incident particles & θ is the angle at which particles are deflected. Integrating the above equation will give us the idea of

the total scattering cross section [49]. Quantum mechanical treatment of scattering deals with the wave functions & square of their moduli; as trajectories (as in Rutherford scattering) make no sense in quantum mechanics. The wavefunction has two parts one an incident part and other the scattered part given as [48];

$$\Psi(\vec{r}) = \Phi_{inc}(\vec{r}) + \Phi_{sc}(\vec{r}) \simeq A \left[e^{i\vec{k}_0 \cdot \vec{r}} + f(\theta, \phi) \frac{e^{i\vec{k} \cdot \vec{r}}}{r} \right]$$

Where $\Phi_{inc}(\vec{r})$ & $\Phi_{sc}(\vec{r})$ are incident and scattered wave functions, A is normalization constant \vec{k}_0 & \vec{k} are incident and scattered wave vectors, $f(\theta, \phi)$ is scattering amplitude. Differential scattering cross section is related to scattering amplitude as [48]

$$\frac{d\sigma}{d\Omega} = \frac{k}{k_0} |f(\theta, \phi)|^2$$

After solving the Schrödinger's equation in a potential V in center of mass frame we get [48];

$$\frac{d\sigma}{d\Omega} = \frac{\mu^2}{4\pi^2 \hbar^4} |\langle \phi | V | \Psi \rangle|^2$$

The methods such as Born approximation (for weak potential) & partial wave analysis can be used for further treatment on differential cross section [50-51].

3.7.2 Neutron scattering.

Neutrons are electrically neutral spin 1/2 particles with an atomic mass of 1.00866491560, the magnetic moment of -1.91304273 nucleon magneton (μ_N) & a lifetime of 885.7 sec [52]. The scattering of neutrons can occur due to either the nucleus or the magnetic moment of the target atoms [53]. The former is referred to as nuclear scattering and the latter is referred to as magnetic scattering [53]. Because of neutron's interaction with the magnetic field, neutron scattering has been an influential tool for the understanding of magnetism [54]. There are different techniques based on neutrons which one can use such as neutron diffraction, small angle neutron scattering, inelastic neutron scattering & quasi-elastic neutron scattering [55]. Neutron diffraction is used for studying crystalline structure of matter [56], inelastic neutron scattering to study vibrational

modes [57], quasi-elastic neutron scattering to study the rotational & translational motion of molecules [58-59], small angle neutron scattering (SANS) to determine sizes & morphologies of materials [60]. SANS is the technique used to characterize the BaTiO₃ nanoparticles (chapter 6).

The atoms can be represented as an ensemble of different scattering lengths b_i (same as scattering amplitude defined earlier) [53]. In general, for an atomic ensemble a new quantity can be defined known as scattering length density (SLD) as [53];

$$\rho = \frac{\sum_i^n b_i}{V}$$

Where V is the volume of target atoms & b_i 's are the scattering lengths of individual atoms in a molecule. Macroscopic scattering cross section then is given as [53];

$$\frac{d\Sigma}{d\Omega}(q) = \frac{1}{V} \left| \int \rho(\vec{r}) e^{i\vec{q}\cdot\vec{r}} d\vec{r} \right|^2$$

Where q is a wave vector given as $q = \frac{4\pi}{\lambda} \sin(\frac{\theta}{2})$, λ is the wavelength of incident radiation & θ is scattering angle.

In an actual experiment, the sample is not suspended in a vacuum, but it is surrounded by external media such as a solvent. So now we have to take into account the role of solvent too that's where the concept of contrast variation emerges. So the actual scattering length density can be given after removing the solvent's contribution as [61];

$$g(\vec{r}) = \rho(\vec{r}) - \rho_s$$

Where ρ_s is scattering length density of the solvent. The more the difference between SLD of particles & solvents more is the contrast, and more clearly we can determine the particle's properties [61]. There are two regimes of scattering one due to Porod & other due to Guinier. Porod's law states that at large q (i.e., larger than the inverse of scattering dimension of the particle) [53,62];

$$I(q) \propto q^{-4}$$

follows. While the low q regime is dealt by Guinier's law which states [53];

$$I(q) \simeq I(0) \exp\left(-\frac{q^2 R_g^2}{3}\right)$$

Where $I(0)$ is forward scattering intensity & R_g is radius of gyration [63].

Form factor $P(q)$ & structure factor $S(q)$ are one of the most important quantities in any scattering experiment. Former refers to scattering by different spatial regions of same particle and latter refers to scattering by different particles [53].

Every particle shape (like sphere, cylinder etc.) has its own form factor. For e.g. spherical particles are identified with form factor [64].

$$P(q) = \frac{3 \sin(qr) - qr \cos(qr)}{qr^3}$$

And when the particles are simple (i.e., no interaction between two particles) structure factor is given as [53];

$$S(q) = 1 + 4\pi N_p \int_0^\infty [g(r) - 1] \frac{\sin(qr)}{qr} r^2 dr$$

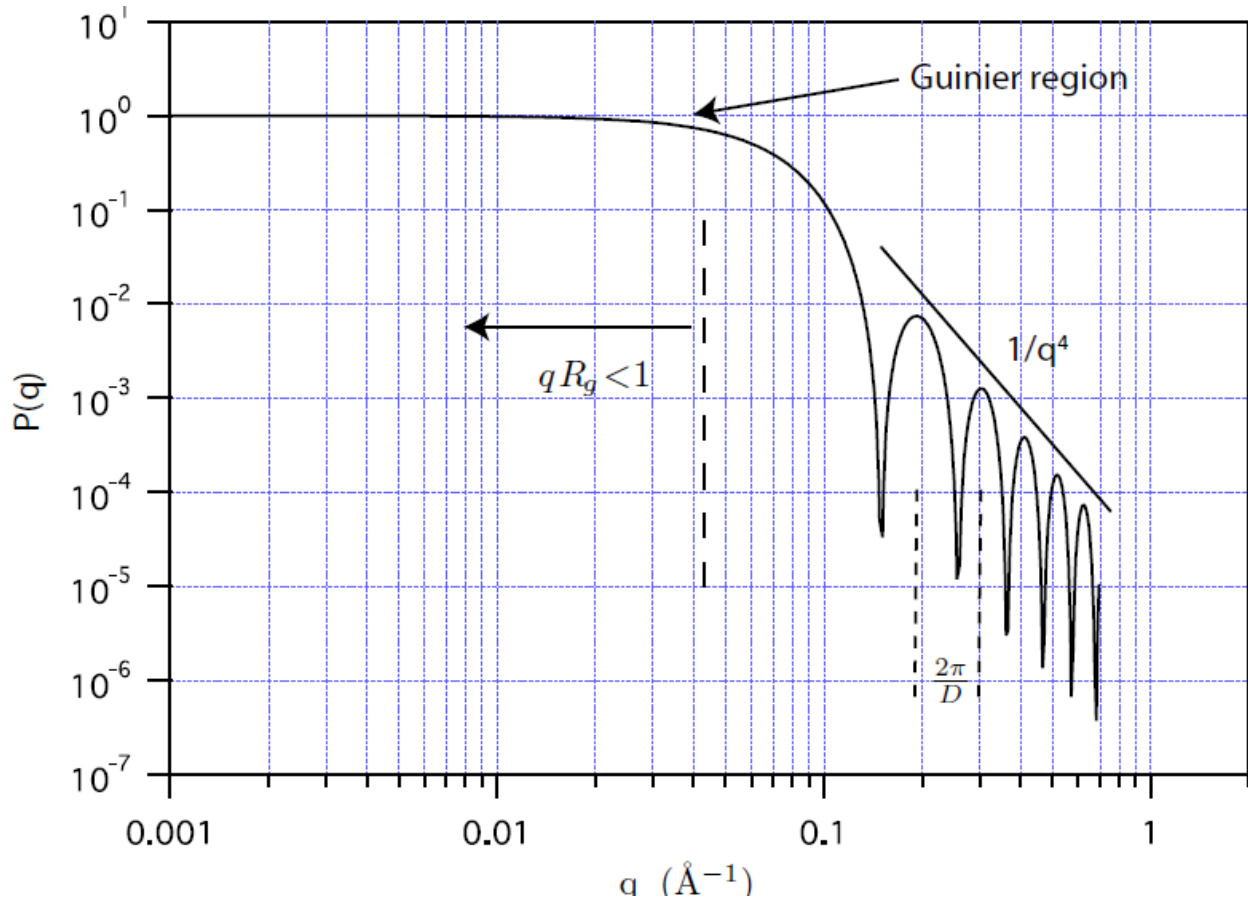


fig 3.11 Plot of $P(q)$ vs q showing guinier & Porod regions.

So in actual analysis, appropriate form factor can be used for fitting to determine the shape of the particles since intensity is directly proportional to both $P(q)$ & $S(q)$ [53];

$$\frac{d\Sigma}{d\Omega}(q) = \frac{N}{V}(\rho - \rho_s)^2 V_p^2 P(q) S(q)$$

Where $V_p = \frac{4}{3}\pi R_g^3$. The fitting of SANS data was done with SASfit software.

3.8 Photoluminescence spectroscopy.

Photoluminescence spectroscopy is a very essential tool to characterize the defects [65]. No crystalline substance is perfect. Every crystalline structure present in nature has defects of some kind or the other that the crystal is to be imperfect is very important by a point of view of

material science and engineering because of the properties they exhibit [65]. The defects have a physical explanation in the view of statistical thermodynamics that due to the increased number of microstates the entropy S is maximized and nature favors the same. The defects present in crystals are mainly of 3 types: point defects, line defects, and surface defects. We are mostly concerned about point and surface defects in this discussion. A point defect in the crystal is a defect due to missing atom from lattice site or extra atom introduced in a lattice site [66]. While the surface defects are the presence of vacancies (atoms missing from their lattice position) or interstitials (atoms acquiring extra lattice positions) on the surface among which the former trap one or more electrons leading to the formation of charged states [67], to probe these defects, the photoluminescence spectroscopy may be a valuable technique. It is a spectroscopic technique involved in the characterization of intrinsic and extrinsic properties of semiconductors and nanostructures, nondestructiveness being an advantage [68].

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Chapter 4: Synthesis of SiO₂ nanoparticles.

4.1 Introduction.

In this chapter, we'll discuss the syntheses of different phases of SiO₂ nanoparticles. The investigation of the ferromagnetism in such nanoparticles is done with the help of SQUID magnetometer. The other characterizations such as P-XRD, photoluminescence spectroscopy, UV visible spectroscopy, fourier transform infrared spectroscopy (FT-IR) & scanning electron microscopy (SEM) are also done to study the origin of ferromagnetism in such nanoparticles.

4.2 Synthesis of SiO₂ nanoparticles of different sizes.

4.2.1 Synthesis of nanoparticles with size~440nm (SiO₂ NP 440).

SiO₂ nanoparticles were synthesized by a method described in chapter 2, i.e., Stober's hydrolysis method [1]. In a typical synthesis procedure [2] two solutions A & B were made. 18.75 ml ethanol was mixed with 1.25 mL tetraethoxysilane (TEOS) to form solution A, and to form solution B 21.5 mL ethanol was mixed with 3.5 mL aqueous ammonia (NH₄OH 25% in water) & 5mL H₂O. Then the solution A was dripped into solution B with a rate lower than 0.05 mL/s with constant stirring. The onset of turbidity was a signature of formation of SiO₂ nanoparticles. The resulting solution was aged for 48h. The solution was then taken for centrifuging and was centrifuged at 10,000 rpm for 5 minutes. The SiO₂ powder thus obtained was washed with ethanol six times and was dried at 60⁰ C overnight.

4.2.2 Synthesis of nanoparticles with size ~200nm (SiO₂ NP 200).

Although the SiO₂ nanoparticles obtained by the previous method showed ferromagnetic nature, the nature is very weak (as will be seen in the discussions of M vs. H data that will follow), so there is a need to synthesize even smaller particles to enhance the magnetic moment. The idea is that smaller the size of nanoparticles, larger will be the specific surface area and in turn magnetic moment [3-6]. In a typical synthesis procedure [7] 1.25 mL TEOS was added in 30mL of EtOH and stirred vigorously for 5-10 minutes then it was allowed to sonicate for 10 minutes after which 1 mL distilled water was added, and the solution was again stirred for 5-10 minutes. Then

it was kept for sonication for 90 minutes, and then 2 mL of NH_4OH was added accompanied by stirring. The turbidity after 3-4 minutes was a clear indication of the formation of nanoparticles. The turbid solution thus obtained was sonicated for 3 hours and aged overnight. After aging, it was centrifuged and washed several times with ethanol and dried for 1-2 days at 60°C in a hot air oven.

4.2.3 Synthesis of nanoparticles of size~100 nm. (SiO_2 NP 100).

To make the nanoparticles of even smaller size another protocol is used which uses sodium silicate as a precursor [8]. In these, the two batches are produced one with a lower volume of solvent let's call it batch 1 (B1), and other with a higher amount of volume let's call it batch 2 (B2). In both syntheses, the volume ratio of ammonia: ethanol was kept the same, i.e., 3:1. In case of synthesis of B1 0.5mL sodium silicate solution(SSS) was dissolved in 7mL of distilled H_2O which was added to the solution of 60 mL ammonia and 20mL ethanol. The final solution became more turbid after aging for 48 hours and then was centrifuged and washed with distilled water and dried at 60°C . For the case of B2 same procedure was implemented except the volume taken was that of 90mL of ammonia and 30mL of ethanol. The reaction process is the same as that of the previous method except that the precursor used is different, i.e., sodium silicate instead of TEOS.

4.2.4 Synthesis of mesoporous nanoparticles (mesoporous SiO_2 NP).

Mesoporous nanoparticles are the ones which are more porous and thus have a high specific surface area (surface area per unit gram) as compared to the conventional nanoparticles (NPs) [9]. These are highly applicable in drug delivery due to the high specific surface area [10]. The aim to prepare such nanoparticles is to enhance the surface area & thus to increase the number of defects present (on the surface) which in turn will increase the magnetic moment. The mesoporous SiO_2 NPs were synthesized as in [11]. In a typical synthesis of mesoporous silica, 32mL of distilled water, 6.5g of ethanol, 5.2g of cetrimonium chloride and 0.1g of diethanolamine (DEA) were mixed and stirred in an oil bath at 60°C for 30-35 minutes. Then 3.65mL TEOS was quickly added dropwise while stirring. The solution turned white due to the formation of amorphous silica and was kept for stirring for 2 hours. The as-prepared turbid solution was centrifuged and washed 3-4 times with distilled water after each run and dried at 60°C in hot oven.

4.2.5 Synthesis of hollow nanoparticles (hollow SiO₂ NP).

Micelles were used as a template to synthesize hollow SiO₂ nanoparticles [12]. 0.08g of cetrimonium bromide (CTAB) was dissolved in 50 mL solution of ammonia. Ammonia: water volume ratio being 0.3. Then 0.5 mL TEOS was added followed by addition of 0.5 mL ammonia solution. The turbid solution thus obtained was stirred for additional 6 hours. After which the suspension was centrifuged and washed with water & ethanol after subsequent runs. The as-obtained powder was kept for drying at 60⁰C in a hot oven.

4.3 Characterization of Nanoparticles.

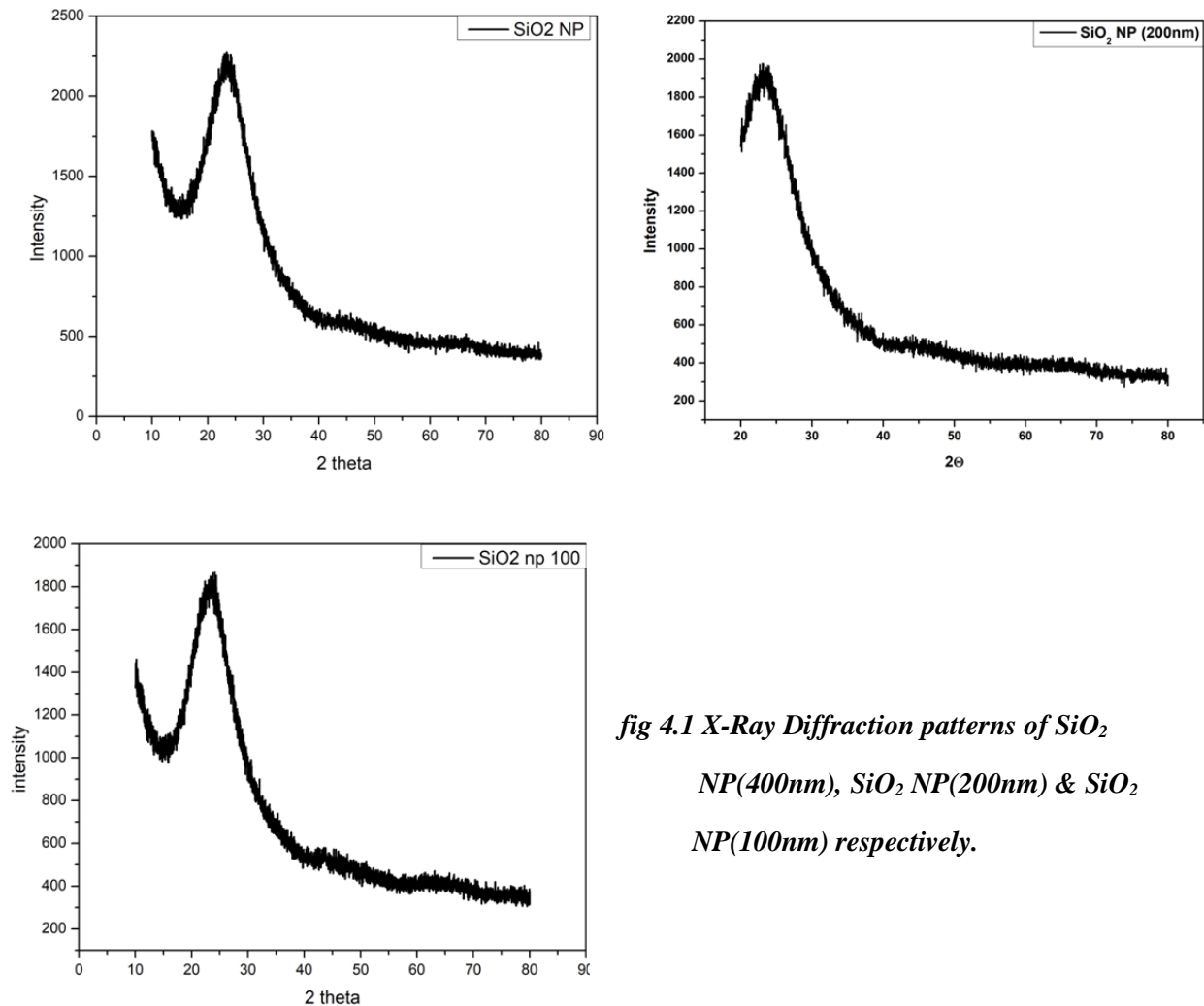


fig 4.1 X-Ray Diffraction patterns of SiO₂ NP(400nm), SiO₂ NP(200nm) & SiO₂ NP(100nm) respectively.

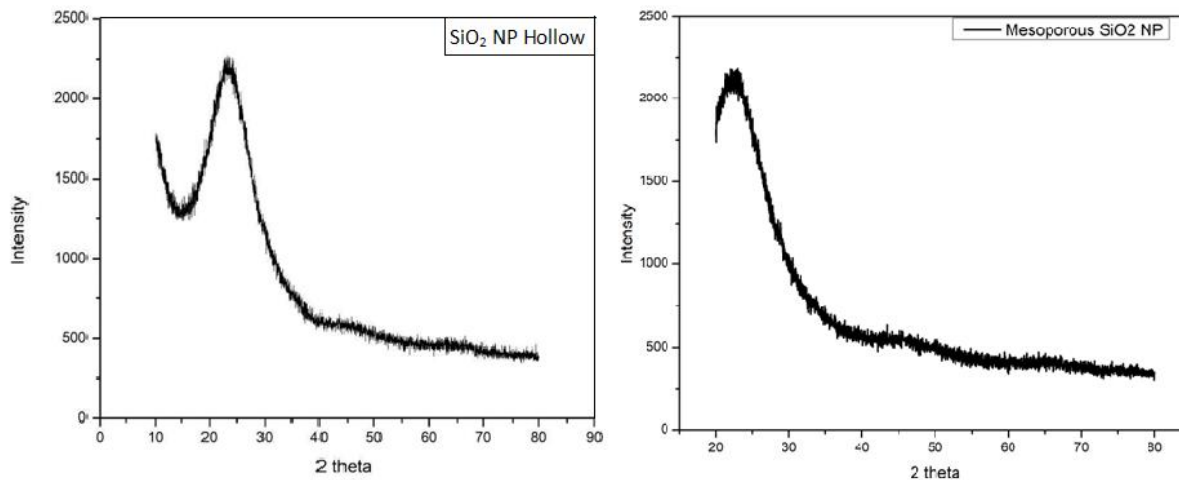


fig 4.2 XRD patterns of both Hollow & Mesoporous SiO₂ Nanoparticles.

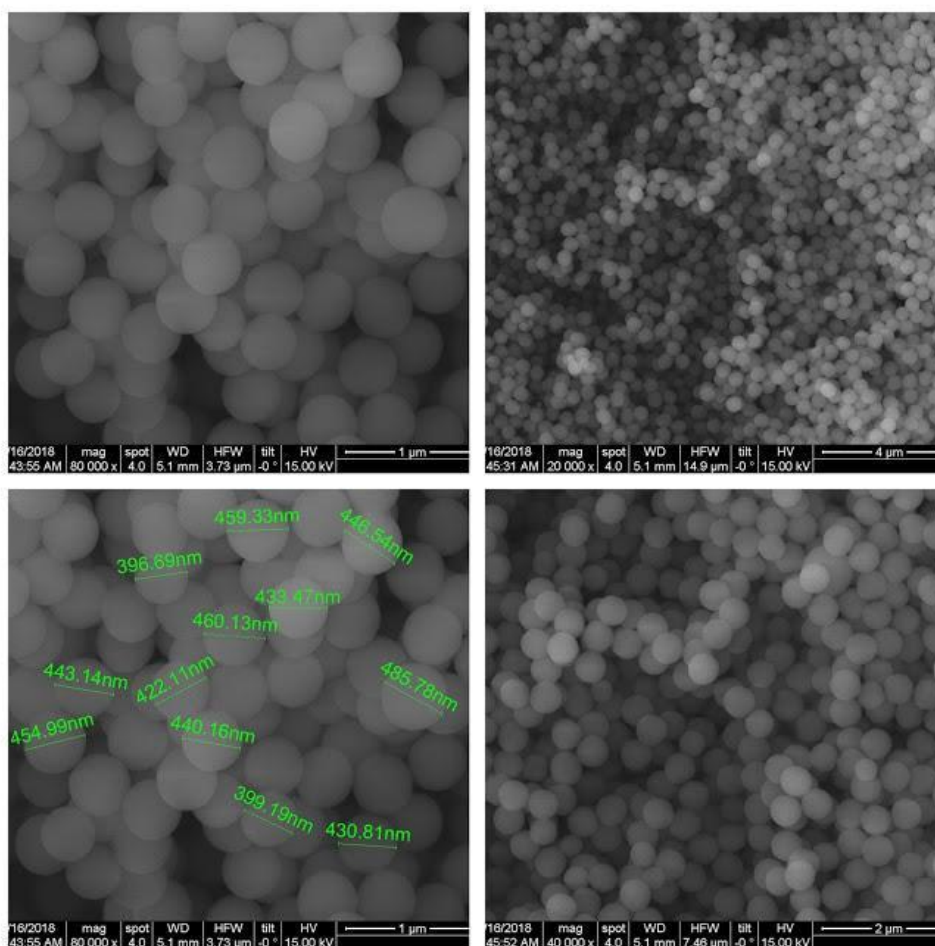


fig 4.3 SEM images of SiO₂ NP(440nm).

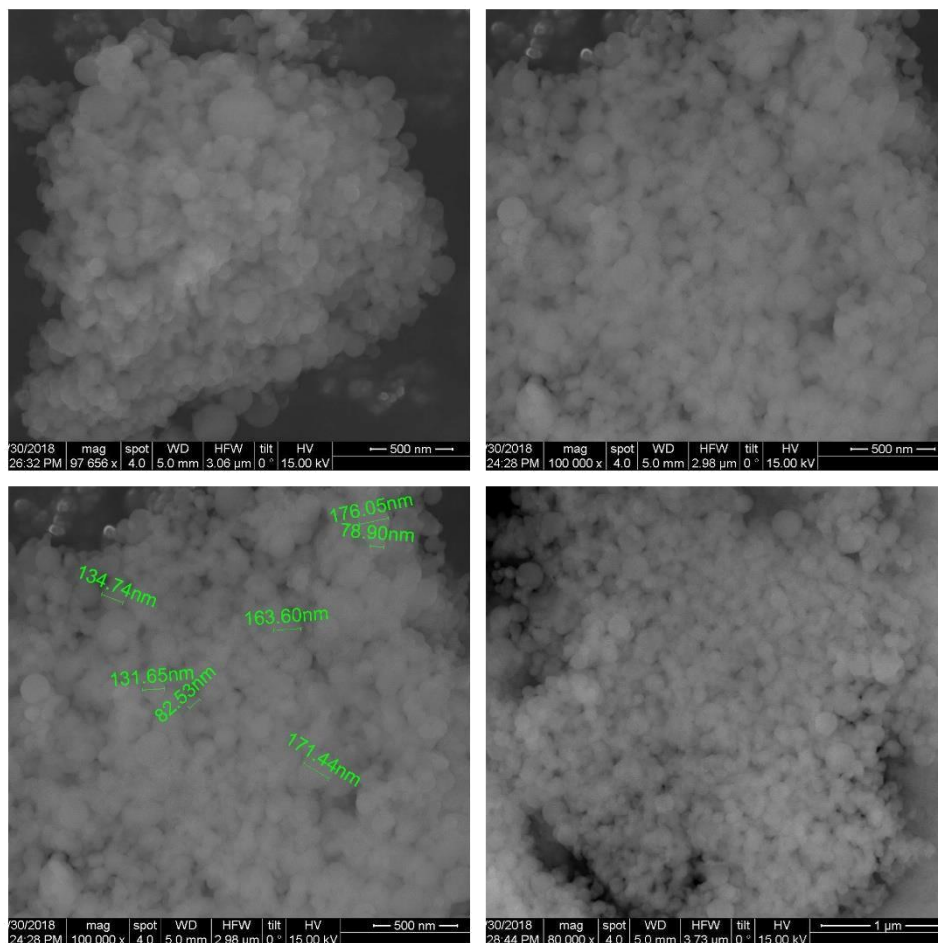


fig 4.4 SEM images of SiO₂ NP (100nm) Batch B1.

The XRD of as-prepared samples were taken with 2θ varying from 10^0 to 80^0 . All the prepared nanoparticles were amorphous as evident from figs 4.1 & 4.2.

The absence of distinct sharp peaks shows the amorphous nature of the nanoparticles. For studying the morphology of as-prepared SiO₂ nanoparticles, the SEM was done. The nanoparticles formed were perfectly spherical, uniform in size and monodispersed in the case of SiO₂ NP (440nm) as is shown in fig 4.3. The SiO₂ NP(200nm) are also uniformly sized, spherical & monodispersed (fig 4.6). The batch B1 of SiO₂ NP (100nm) as shown in fig 4.4 doesn't show uniform spherical & monodispersed nanoparticles. But batch B2 shows much-improved uniformity in size (fig 4.5).

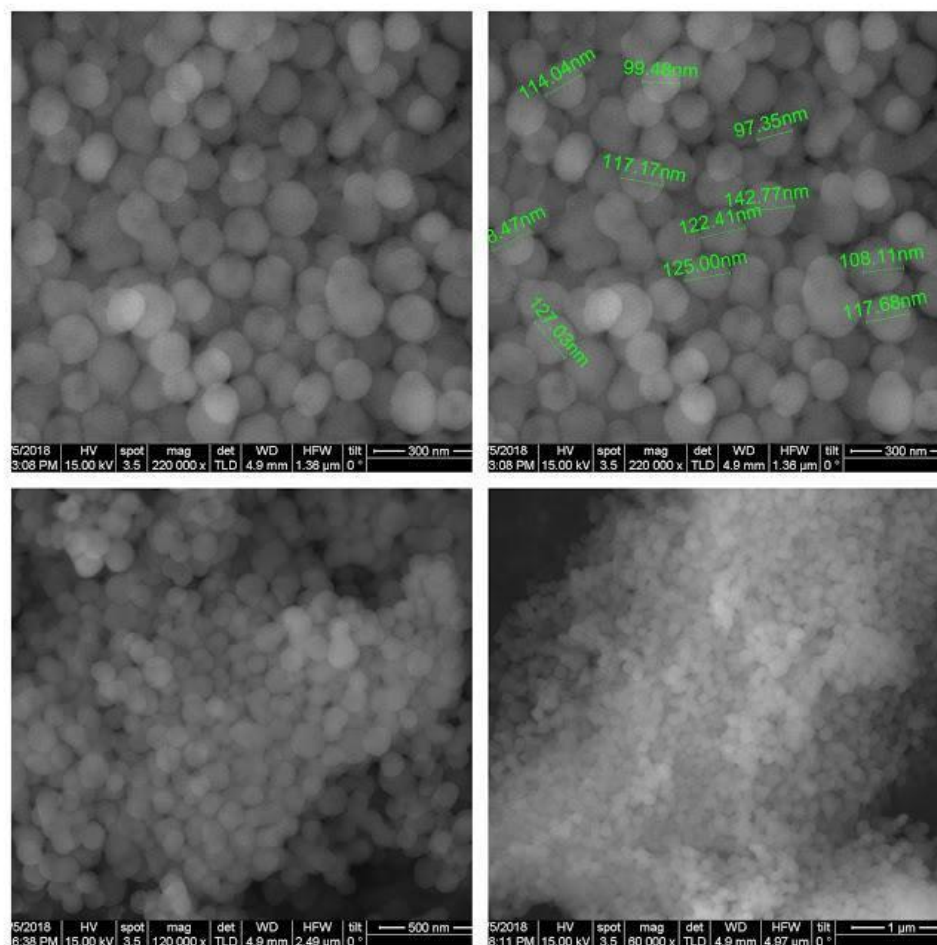


fig 4.5 SEM images of SiO₂ NP (100nm) Batch B2.

The magnetic measurements were done with SQUID magnetometer. M vs. H data was taken for the nanoparticles. The magnetic data shown in fig 4.7 was taken at 300K temperature. The data of both types of NPs clearly show the room temperature ferromagnetism. Both are soft magnetic. The SiO₂ NP (100nm) shows about three times more saturation than that of SiO₂ NP (440nm) which may be due to higher surface area; which is in agreement with previous work by Prof. Sundaresan et al. on other compounds [3-6].

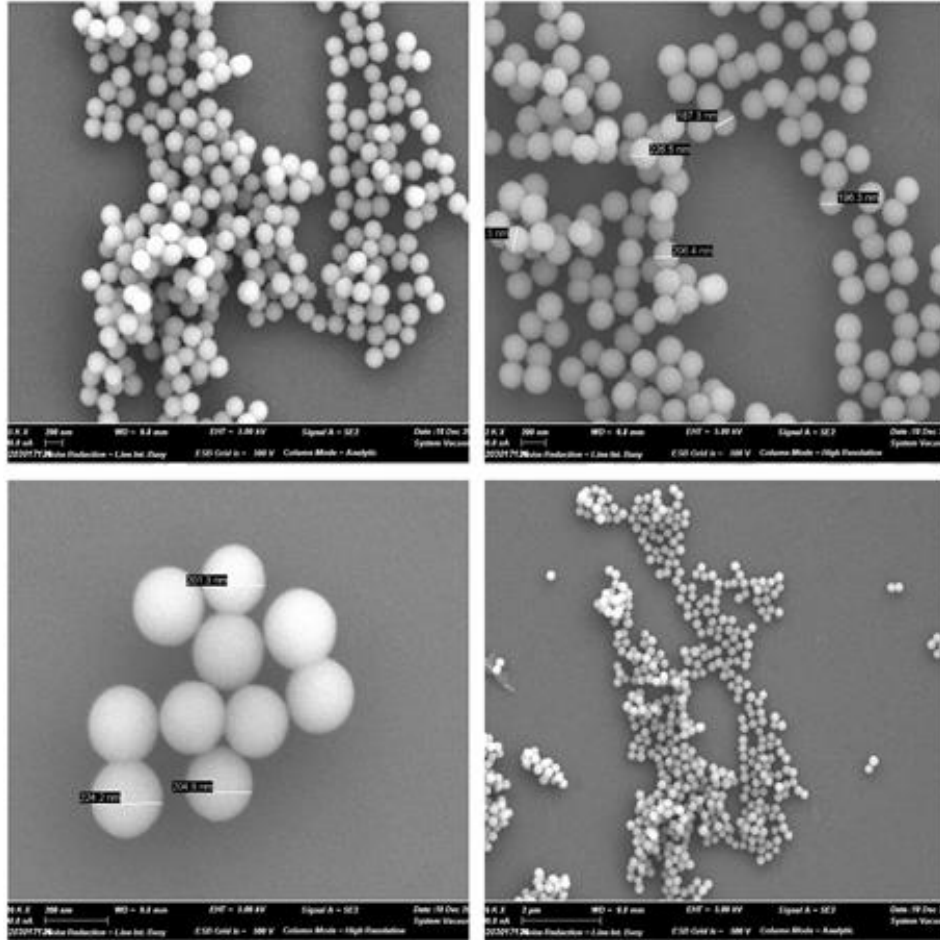


fig 4.6 SEM of SiO₂ NP (200nm).

The more saturation may be due to the higher number of concentration (per gram) of oxygen vacancies due to the higher surface area. In the case of mesoporous SiO₂ nanoparticles magnetic moment is too low as can be seen in fig 4.8. The result is rather surprising and against general expectation, because the moment anticipated was very high because of the high BET surface area of about 722 m²/g [11]. While the hollow nano spheres show much higher moment as compared to SiO₂NP(440nm) as can be seen in fig 4.9 but yet lower than that of SiO₂ NP(100nm) still it is a relatively good result for estimated size of 250-500nm [12].

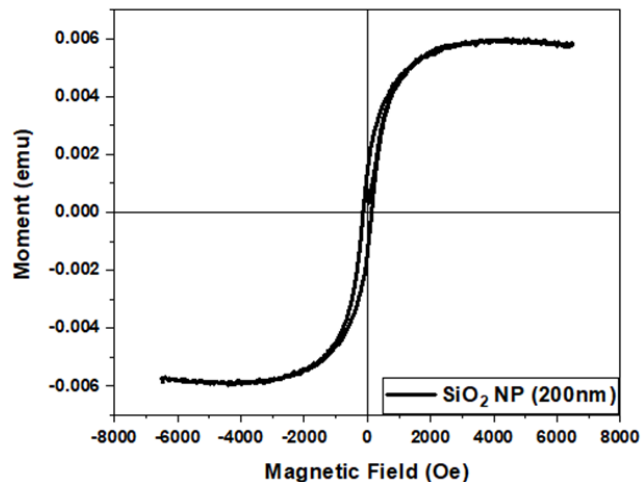
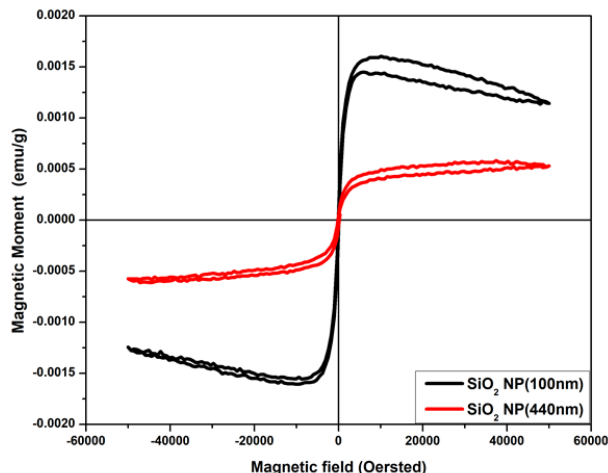


fig 4.7 *M vs H* data of SiO₂ NP (440nm) & SiO₂ NP (100nm). fig 4.33 *M vs. H* of SiO₂ NP (200nm).

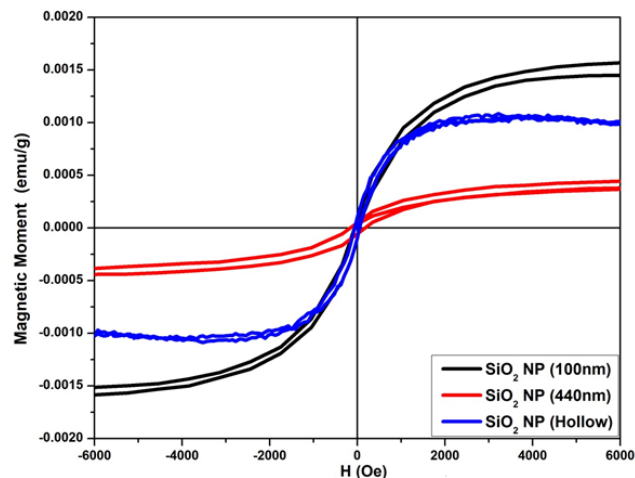
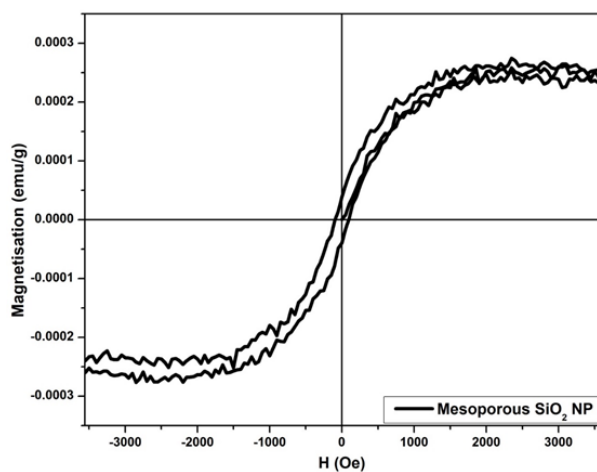


fig 4.8 *M vs. H* data of Mesoporous SiO₂ NP.

fig 4.9 *M vs. H* data of Hollow SiO₂ NPs

compared to both 100nm & 440 nm sized nanoparticles.

SiO₂ NP (200nm) shows nearly six times higher magnetic moment than SiO₂ NP (100nm) (fig 4.33) which is surprising given its relatively bigger size. It may be due to the ultrasonication involved in its synthesis. After annealing, the magnetic moment goes down to half of its initial value.

UV visible studies were also done to find the band gap of as-prepared nanoparticles. Fig 4.10 & fig 4.11 shows the UV visible data of as prepared SiO₂ nanoparticles. The band gap was calculated by using the Tauc plot method.

4.3.1 Tauc Plots.

The valence electrons of atoms which are lower in energy are called valence band electrons, while the ones which are higher in energy and mostly responsible for conduction are conduction band electrons [14]. On the incidence of light with a particular frequency, electrons undergo a transition from valence band to conduction band [15]. The method usually used in the determination of band gap of materials from UV Visible spectra is known as Tauc plot.

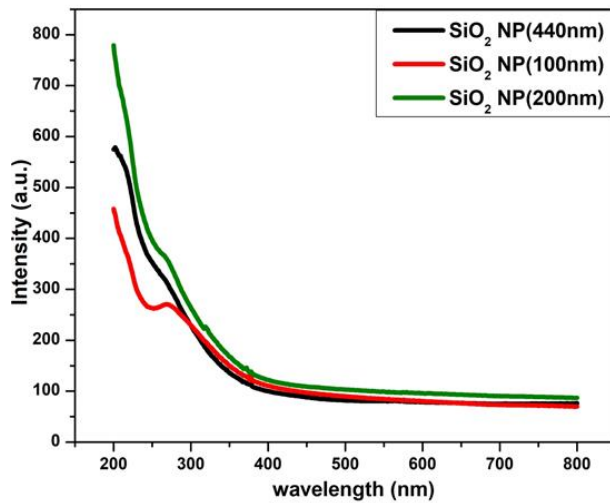


fig 4.10 Size dependence on UV visible spectra of SiO₂ Nanoparticles.

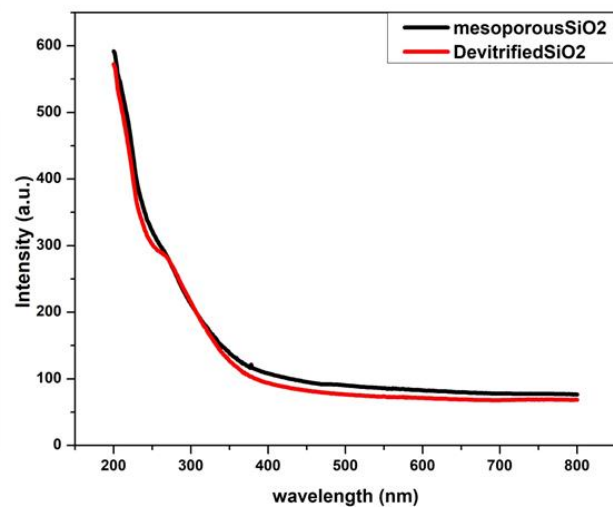


fig 4.11 UV Visible spectra of Mesoporous & Hollow SiO₂ NP.

To determine the band gap; understanding of Kubelka-Munk (KM) theory is essential. It assumes that diffuse reflection is a result of both scattering and absorption as evident from KM equation [13];

$$F(R_{\infty}) = \frac{(1 - R_{\infty})^2}{2R_{\infty}} = \frac{K}{S}$$

Where $F(R_{\infty})$ is KM function, R_{∞} is absolute Remittance, K is absorption coefficient & S is

the scattering coefficient. After knowing $F(R_{\infty})$ band gap can be determined by the equation [16],

$$(F(R_{\infty})hv)^2 = A^2(hv - E_g)$$

Where hv is energy, E_g is band gap.

So the first thing done was to convert 'x' axis of wavelength (nm) to energy(eV). The y-axis was changed from R to $F(R_{\infty})$ & the graph was plotted. Then the extrapolation was done as shown in figs 4.12,4.13 & 4.14. The intercept on the x-axis is the bandgap of the material. The results of the band gap calculations are shown in Table 4.1.

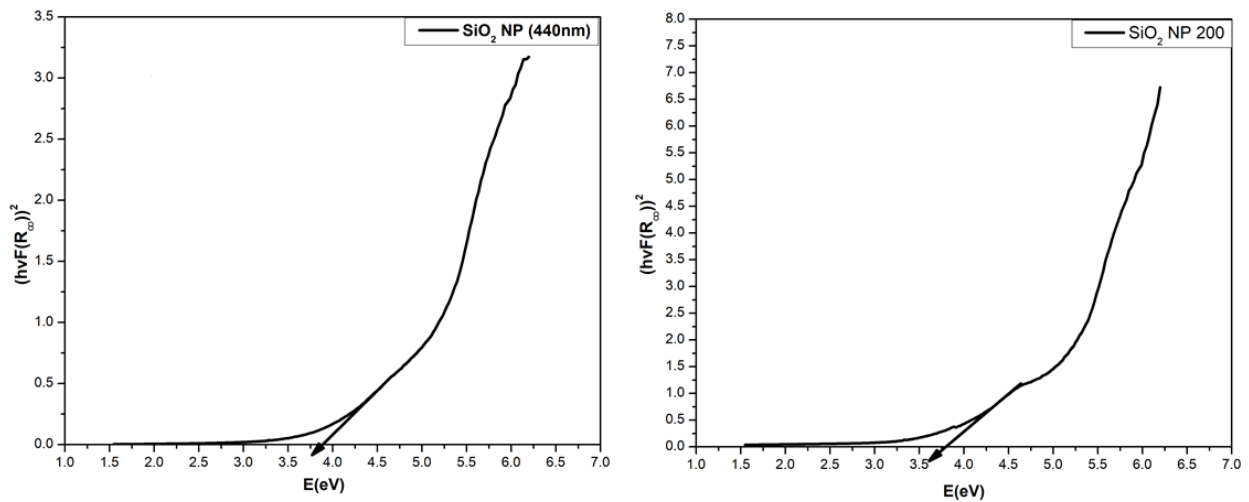


fig 4.12 Tauc plots of SiO₂ NP (440nm) & SiO₂ NP (200nm).

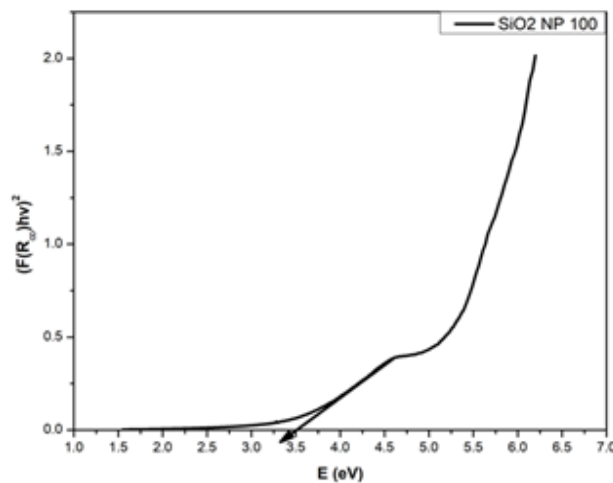


fig 4.13 Tauc plot of SiO₂ NP (100nm).

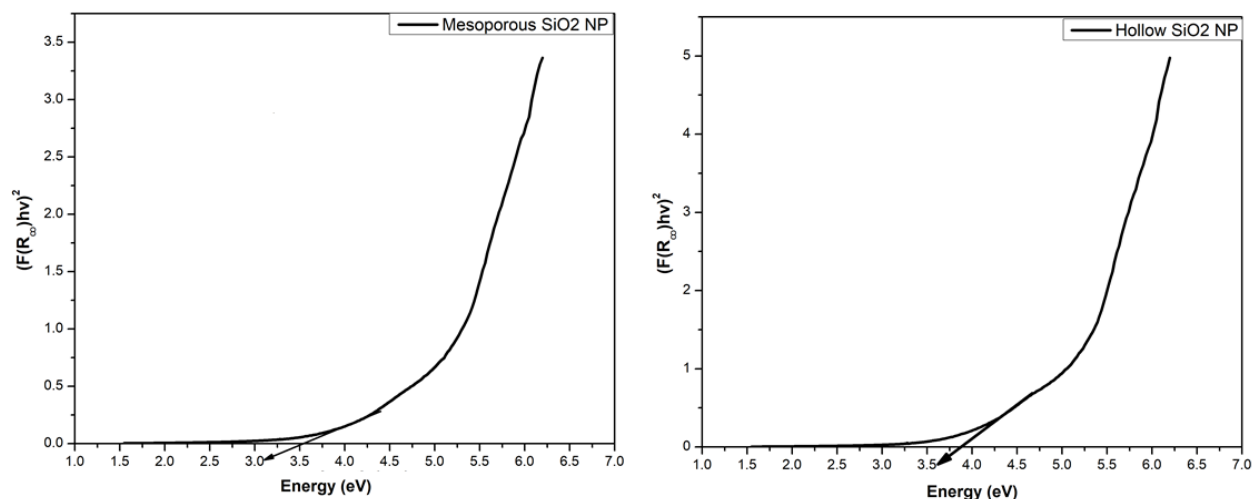


fig 4.14 Tauc Plots of Hollow & Mesoporous SiO₂ NP.

As can be seen from Table 4.1 band gap of nanoparticles decrease as the particle size decrease. The Mesoporous & Hollow nanoparticles both show a band gap of 3.56 eV & 3.88 eV respectively.

Sample Name	Band Gap (eV)
SiO ₂ NP (440nm)	3.88
SiO ₂ NP (200nm)	3.79
SiO ₂ NP (100nm)	3.54
Mesoporous SiO ₂ NP	3.56
Hollow SiO ₂ NP	3.88

Table 4.1 Band gap of as prepared SiO₂ Nanoparticles.

The IR spectra of the nanoparticles were also taken to see the types of chemical bonds present in the nanoparticles. Figs 4.15 & 4.16 both show the FT-IR (Fourier transform infrared) spectra of the nanoparticles. The bond characterization is done as follows;

4.3.2 Bond Characterization.

4.3.2.1 Si-OH stretching vibrations.

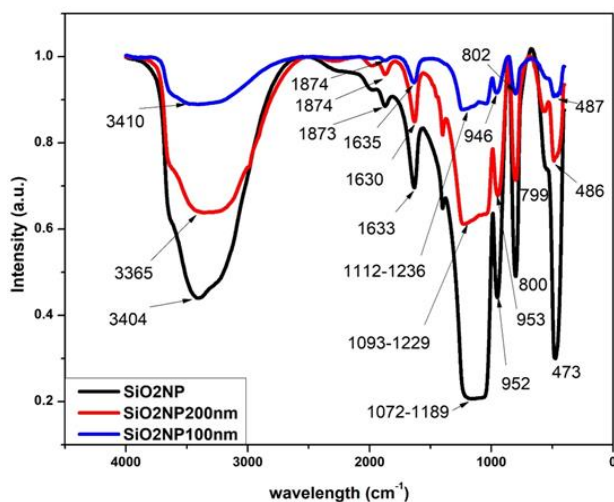


fig 4.15 Infra-red spectra of different sizes of SiO₂ Nanoparticles.

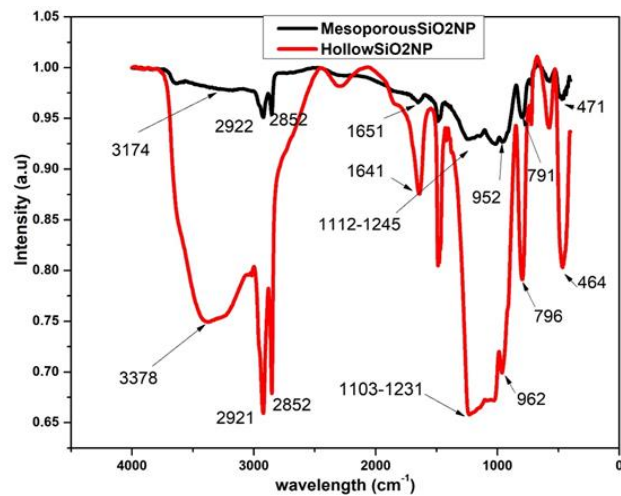


fig 4.16 Infra-red spectra of Hollow & Mesoporous SiO₂ Nanoparticles.

The characteristic band of SiO₂ NP (440nm), SiO₂ NP (200nm) & SiO₂ NP (100nm) at 3404, 3365 & 3410 cm⁻¹ respectively is speculated to be from Si-OH stretching mode [17-19]. The same band occurs at 3174 & 3378 cm⁻¹ for mesoporous & Hollow nanoparticles respectively.

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	3404
SiO ₂ NP (200nm)	3365
SiO ₂ NP (100nm)	3410
Mesoporous SiO ₂ NP	3174

Hollow SiO₂ NP	3378
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Table 4.2 Peak positions of Si-OH stretching.

But for mesoporous nanoparticles, it is nearly nonexistent. The band origin may be due to H₂O molecules adsorbed on nanoparticle surface [20]. But some contribution to this band can also be from bridging hydrogen between the water molecules given its broadness [21] which indicates the presence of moisture which may be absorbed by KBr while making pellets. The number of silanols per unit gram decrease as the nanoparticle size of the sample decreases. And in the case of mesoporous nanoparticles, there may be negligible numbers of silanol groups.

4.3.2.2 O-H Bending.

The existence of Si-OH stretching can be confirmed by the O-H bending bands. These bands can be located as shown in table 4.3. The band is due to water molecule [22], which confirms the existence of Si-OH on the surface.

Sample Name	Peak Position (cm⁻¹)
SiO₂ NP (440nm)	1633
SiO₂ NP (200nm)	1630
SiO₂ NP (100nm)	1635
Mesoporous SiO₂ NP	1651
Hollow SiO₂ NP	1641

Table 4.3 Peak positions for OH bending.

4.3.2.3 Si-O-Si Asymmetric vibration.

The Si-O-Si Asymmetric modes are as assigned in Table 4.4 [23-25].

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	1072-1189
SiO ₂ NP (200nm)	1093-1229
SiO ₂ NP (100nm)	1112-1236
Mesoporous SiO ₂ NP	1112-1245
Hollow SiO ₂ NP	1103-1231

Table 4.4 Peak positions of Si-O-Si Asymmetric stretch.

4.3.2.4 Si-OH bending vibration.

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	952
SiO ₂ NP (200nm)	946
SiO ₂ NP (100nm)	953
Mesoporous SiO ₂ NP	952
Hollow SiO ₂ NP	962

Table 4.5 Peaks assigned for Si-OH bending vibrations.

The Si-O-H bending vibrations are assigned as can be seen in table 4.5 [26].

4.3.2.5 Si-O Bending Vibrations.

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	800
SiO ₂ NP (200nm)	799
SiO ₂ NP (100nm)	802
Mesoporous SiO ₂ NP	791
Hollow SiO ₂ NP	796

Table 4.6 Peaks assigned for Si-O bending vibrations.

The peak positions of Si-O bending vibrations are shown in Table 4.6[27-29].

4.3.2.6 Si-O Rocking Vibrations.

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	473
SiO ₂ NP (200nm)	486
SiO ₂ NP (100nm)	487
Mesoporous SiO ₂ NP	471
Hollow SiO ₂ NP	464

Table 4.7 Peaks assigned for Si-O Rocking Vibrations.

The Si-O rocking vibrations are shown in table 4.7 [30-32].

4.3.2.7 –O-CH₃ stretching Vibrations.

Sample Name	Peak Position (cm ⁻¹)
Mesoporous SiO ₂ NP	2922
Hollow SiO ₂ NP	2921

Table 4.8 Peaks assigned for -O-CH₃ stretching vibrations.

The -O-CH₃ stretching vibrations are shown in table 4.8 [33,34]. The –O-CH₃ bands are only observed for mesoporous & hollow SiO₂ NP. The presence of -CH₃ may be due to involvement of ethanol on the surface of nanoparticles.

4.3.2.8 Symmetric C=O stretching.

Sample Name	Peak Position (cm ⁻¹)
SiO ₂ NP (440nm)	1873
SiO ₂ NP (200nm)	1874
SiO ₂ NP (100nm)	1874

Table 4.9 Peaks assigned for symmetric C=O stretching.

The symmetric C=O stretching is only observed for samples depicted in table 4.9. Their peak positions are also shown [35-37].

4.3.2.9 Symmetric CH₂ stretching.

Sample Name	Peak Position (cm ⁻¹)
Mesoporous SiO ₂ NP	2852
Hollow SiO ₂ NP	2852

Table 4.10 Peaks for symmetric CH₂ stretching.

Due to the presence of ethanol as solvent, there may be some aliphatic hydrocarbon moieties present which give rise to symmetric CH₂ stretching as shown in table 4.10 [38-40].

The photoluminescence (PL) spectra of as-prepared nanoparticles of sizes 440nm, 100nm, & hollow SiO₂ NP is as shown in fig 4.17. The excitation wavelength was 325 nm, and emission was observed in the range of 340nm to 575nm of wavelength. The two sharp peaks which are present in both the spectra are not to be considered because those have originated from the sample holder & are very tough to eliminate. In fig 4.18 The Photoluminescence of SiO₂ NP (200nm) is shown in the range 300-500nm, excitation wavelength being 290nm. Its PL intensity is not to be compared with other nanoparticles since it is measured from a different instrument. The mesoporous SiO₂ shows a single broad peak centered at around 388 nm which is almost three times more than SiO₂ NP (440nm), SiO₂ NP (100nm) & Hollow SiO₂ NP (fig 4.19).

In the semiconductors, the anion vacancy or the interstitial can trap electrons to give rise to what is known as excitons [41]. When the light with equal or higher energy is incident on such compounds the electrons trapped in the defects get excited, and upon relaxation they come to ground state only to radiate energy in the form of light or phonons, the former refers to the photoluminescence if the relaxation doesn't last for more than ten seconds. Otherwise, it comes under the category of Phosphorescence. The bands observed for SiO₂ NP (440nm) & SiO₂ NP (100nm) at 431 nm translate to the energy of 2.87 eV.

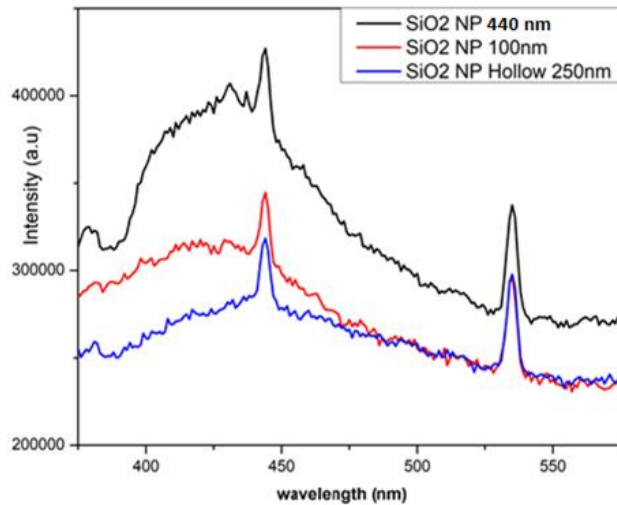


fig 4.17 Photoluminescence of SiO₂ NP (440nm), SiO₂ NP (100nm) & Hollow SiO₂ NP excited at 325nm wavelength.

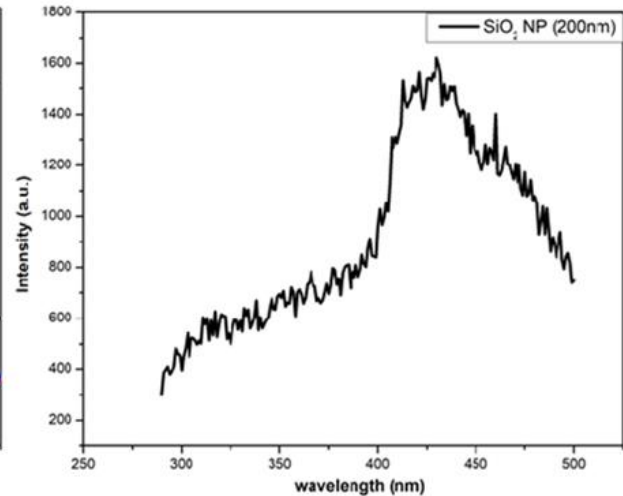


fig 4.18 Photoluminescence of SiO₂ NP (200nm) excited at 290nm.

In the case of SiO₂ NP (200nm), the band is evident at 430nm, i.e., 2.88 eV. The band observed in all cases is due to the excitons and oxygen deficiency E'γ centers [42-45]. For the mesoporous nanoparticles, the band is shifted to 3.1eV, i.e., to higher energy.

4.4 Effect of annealing.

SiO₂ NP (440nm), SiO₂ NP (100nm) & SiO₂ NP (200nm) were annealed at 500⁰ C in hot air oven. The effect of annealing on magnetic properties was studied by SQUID, PL, UV visible spectroscopy & FT-IR.

4.5 Characterization of nanoparticles after annealing.

UV visible spectra of annealed nanoparticles were studied to investigate the effect on the bandgap. Fig 4.20 depicts the UV visible spectra of SiO₂ nanoparticles before and after annealing. The band gap was again calculated using the Tauc plot method.

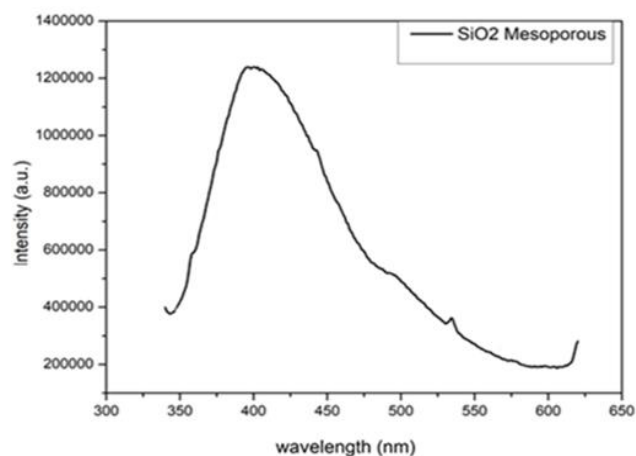


fig 4.19 PL of Mesoporous SiO₂ NP excited at 325 nm showing a broad peak around 388 nm.

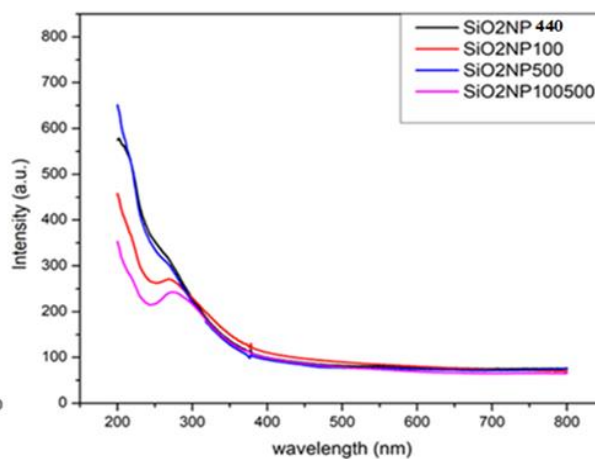


fig 4.20 UV Visible spectra of SiO₂ NPs before and after annealing.

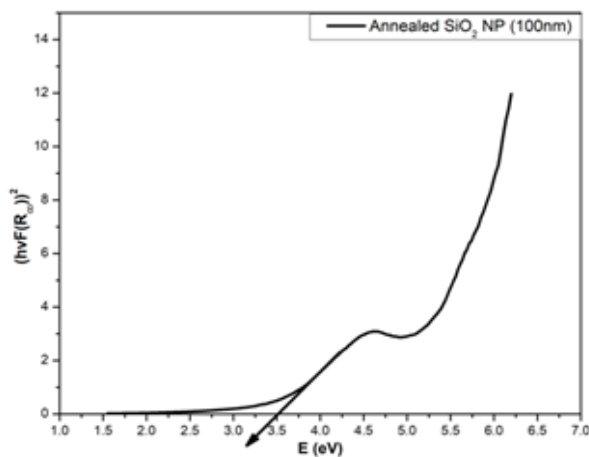
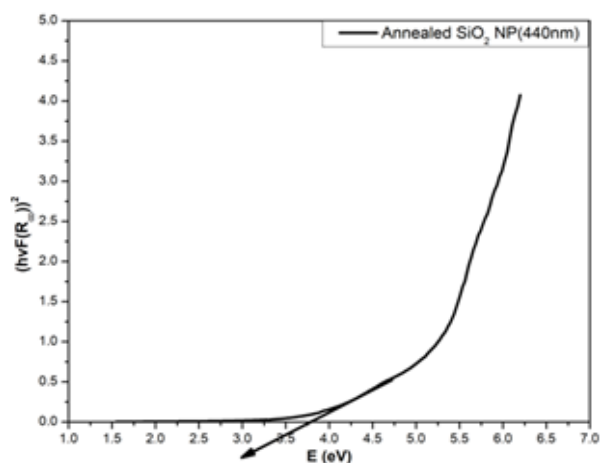


fig 4.21 Tauc plots of annealed SiO₂ NP(440nm) & SiO₂ NP (100nm).

Sample Name	Band gap before annealing	Band gap after annealing
SiO ₂ NP (440nm)	3.88	3.79
SiO ₂ NP (100nm)	3.54	3.52

Table 4.11 Band gaps before and after annealing.

The band gap has decreased after annealing as can be seen in Table 4.11.

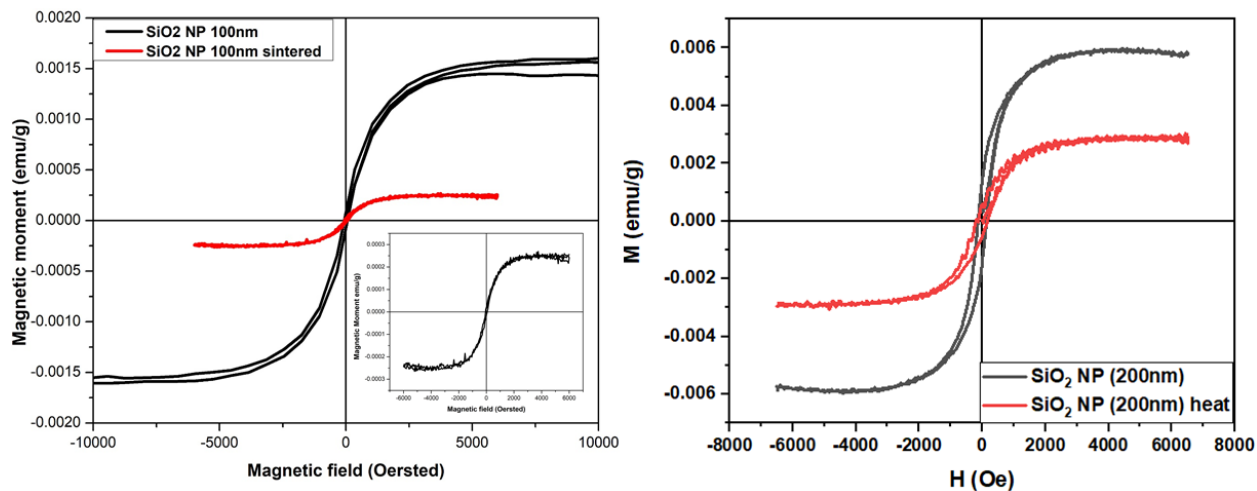


fig 4.22 M vs. H of SiO₂ NP (100nm) before and after annealing (left). The inset shows the magnified image of M vs. H of annealed SiO₂ NP. While the right figure is that of M vs. H of SiO₂ NP (200nm) before and after annealing.

As can be seen from fig 4.22 after annealing of SiO₂ NP (100nm) the saturation magnetism went down from 0.00175 emu/g to 0.00025 emu/g i.e., the magnetism was decreased by a staggering seven times. While in the case of SiO₂ NP (200nm) the magnetic moment has nearly halved after annealing.

The FT-IR of annealed SiO₂ NP (440nm) shows the presence of only four tiny peaks as compared to the original.

SiO ₂ NP (440nm)	Peak Position (cm ⁻¹) (before annealing)	Peak Position (cm ⁻¹) (after annealing)
Si-O-Si asymmetric vibration	1072-1189	1248
Si-OH bending vibration	952	1015
Si-O Bending	800	805
Si-O Rocking	473	476

Table 4.12 Peak positions SiO₂ NP (440nm) before and after annealing.

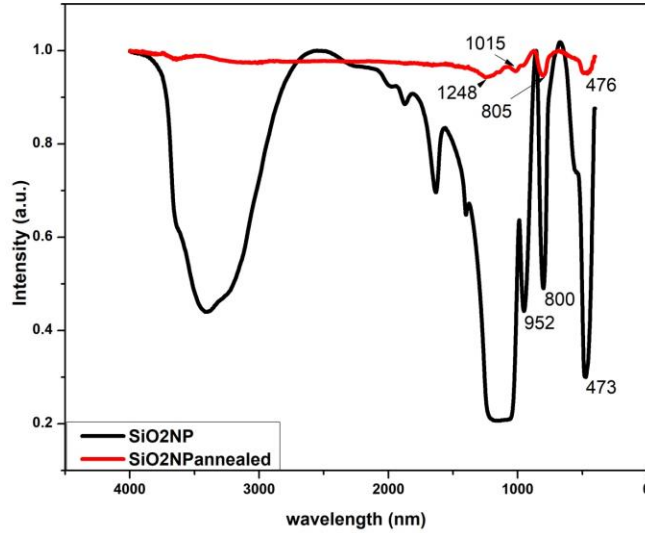
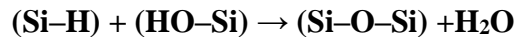


fig 4.23 FT-IR of SiO₂ NP (440nm) before and after annealing.

Only four IR peaks are detected after annealing that too with much lesser intensity as can be seen in fig 4.23. All the peaks shifted to higher energies after annealing which indicates higher energy transitions have occurred after heat treatment. There is a reduction of oxygen functionalities after annealing [46]. The Si-OH (Silanol) concentration has gone down because of the dehydroxylation process as suggested in the equation below [47].



So because of the higher temperature, the water in Silanols evaporates. The high energy shift of Si-O-Si asymmetric vibration shows that the oxygen vacancies are indeed filled out [48,67]. This may have little bit reflected in Si-O bending & rocking mode also.

The Photoluminescence (PL) was done by exciting at 325 nm wavelength.

As can be seen in fig 4.24 & fig 4.25 the PL of SiO₂ NP (440nm) & SiO₂ NP (100nm) after annealing shows lower intensity than the original which indicates decreased oxygen vacancies.

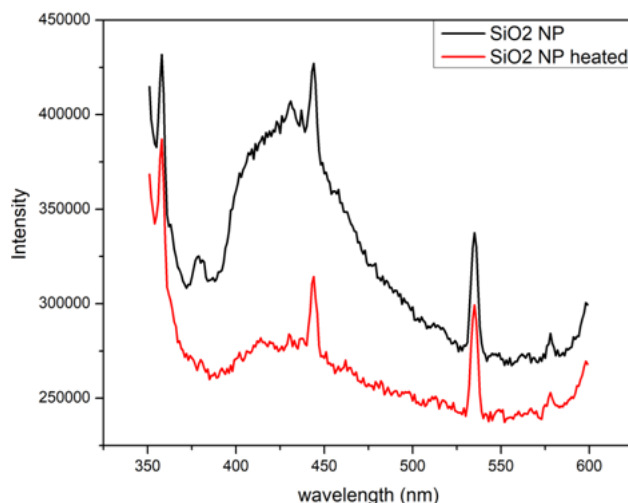


fig 4.24 PL images of SiO₂ NP (440nm) before and after annealing.

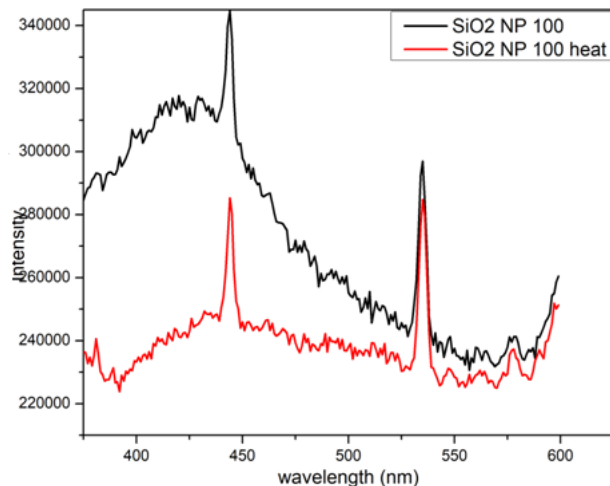


fig 4.25 PL of SiO₂ NP (100nm) before and after annealing

4.6 Results & Discussions.

We have successfully synthesized SiO₂ nanoparticles of sizes 440 nm, 200 nm & 100 nm. We also synthesized hollow & mesoporous SiO₂ nanoparticles. The idea behind synthesizing mesoporous and hollow nanoparticles was to increase the surface area and thus the magnetic moment [3-6].

We've thus investigated the dependence of the magnetism of SiO₂ nanoparticles on the size and temperature treatment. The characterizations such as XRD, SEM, photoluminescence, FT-IR, Diffuse reflectance UV was done to understand the mechanism. As the size of the nanoparticles decreased, the magnetism increased. We know that smaller nanoparticles have a bigger surface area so, in turn, more surface oxygen vacancies so we speculate that the origin of magnetism may be due to oxygen vacancies. After annealing of SiO₂ nanoparticles, the photoluminescence went down and the magnetic moment too went down in case of SiO₂ NP (100nm); so we may conclude that the oxygen vacancies is the reason for the magnetism. The much higher magnetic moment of SiO₂ NP (200nm) (i.e., six times than SiO₂ NP (100nm)) may be because of the synthesis procedure which was involved in its synthesis. About 3 hours and 90 minutes of ultrasonication during synthesis may have introduced defects and oxygen vacancies in the

nanoparticles as seen in its PL spectra. The band gap has decreased with the decrease in nanoparticles size which contradicts consensus [66-67].

4.6.1 Why SiO₂ NP (200nm) has higher magnetic moment than SiO₂ NP (100nm) which has more specific surface area?

The report from Jorge Mejia et al. [68] (which is the study of effect of ultrasonication on the surface composition of nanoparticles in contrast to simple stirring) has shown that the sonication indeed creates C deficiencies on the surface of TiC nanoparticles (XPS studies). While their TEM results also showed the increase in surface roughness in both SiC & TiC nanoparticles after the sonication. Similar must have happened in the case of SiO₂ NP (200nm). The oxygen deficiencies might have been induced on the surface of the SiO₂ nanoparticles due to ultrasonication.

The number of oxygen vacancies in SiO₂ NP (200nm) are so high that they have outnumbered the oxygen vacancies present in SiO₂ NP (100nm) which has more specific surface area than SiO₂ NP (200nm).

4.7 Devitrification & fluorination.

The only ways to enhance the magnetic properties until now are to lower the size or to induce oxygen vacancies by doping (discussed in the next chapter) with other non-magnetic materials, but there can be other ways to enhance magnetism (rather than relying on oxygen vacancies) which are to utilize pairing interactions between the dopant atoms. It is already known that doping with C or N makes group II oxides which are usually nonmagnetic, ferromagnetic [49-50]. The DFT calculations demonstrate the increase in magnetism is due to pairing interactions of dopants itself [51]. We also can try to induce such pairing interactions in SiO₂ nanoparticles but the theoretical calculations done in [51] are only applicable to crystalline solids. Hence we can try to crystallize the amorphous SiO₂ nanoparticles. The process of such crystallization is called devitrification [52]. To crystallize the amorphous SiO₂, we have to break the bridging bonds so SiO₄ molecules can align in their preferred crystalline orientation. So the as-prepared mesoporous SiO₂ nanoparticles were kept in 1M Sr(NO₃)₂ solution for 2-3 hours and then was calcined at 750^oC only to obtain crystallized SiO₂ nanoparticles as can be seen in fig 4.27 [53].

As already stated in chapter 1; studies by Gregory Kopnov et al. have shown that the etching of the silicon wafers by HF and KOH have shown relatively large magnetic moment. Their studies suggest that the magnetism is co-related to surface roughness [54]. But the hole-bridge structures which were obtained by etching the already sintered Nano spherical array with HF show weak ferromagnetism [2]. It is speculated that HF-etching may have contributed to weak ferromagnetism in these Nano spherical arrays [2]. But instead of HF, we have used NH_4F . 0.0067 g of NH_4F was dissolved in H_2O , and 0.03 g of mesoporous SiO_2 sample was added to it. We know on an average 4.9 OH groups per nm^2 exist on the surface of SiO_2 nanoparticles [55]. We also know that BET surface area of as-prepared mesoporous SiO_2 nanoparticles is $722 \text{ m}^2/\text{g}$. So from this information, we have calculated how much OH groups per gram are present on nanoparticle surface and in turn how much moles NH_4F is required to fluorinate the surface.

4.7.1 Characterization of devitrified and fluorinated SiO_2 nanoparticles.

The X-ray diffraction pattern after devitrification shows the onset of crystallinity as can be seen in fig 4.27. While fluorinated SiO_2 nanoparticles show crystalline nature (fig 4.26). The peaks are of $(\text{NH}_4)_2\text{SiF}_6$. UV visible spectra are shown in fig 4.28. The band gap of mesoporous SiO_2 nanoparticles increased from 3.56 eV to 3.78 eV & 3.89 eV for devitrified & fluorinated nanoparticles respectively.

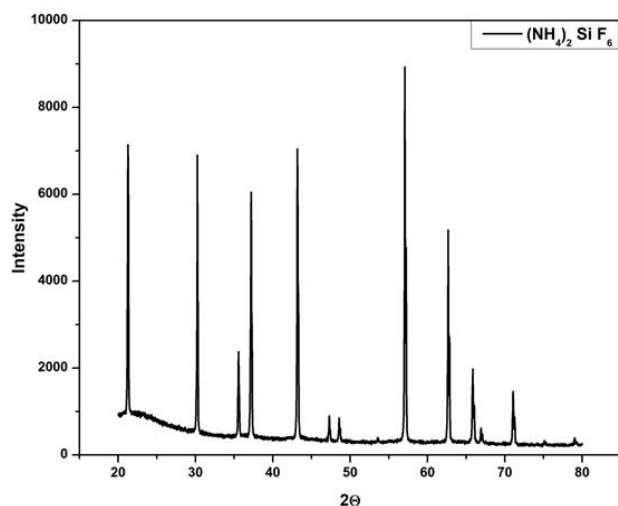


fig 4.26 XRD after fluorination. $(\text{NH}_4)_2\text{SiF}_6$ has formed.

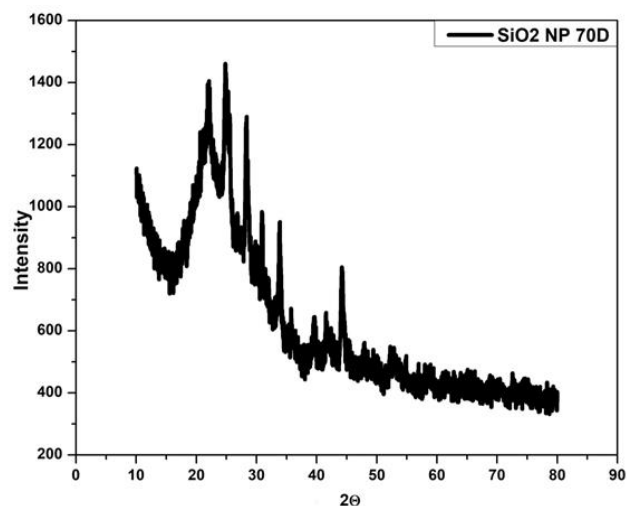


fig 4.27 XRD pattern shows onset of crystallinity although amorphous phase is still present.

The FT-IR spectra of both the devitrified and fluorinated SiO₂ nanoparticles as compared to mesoporous nanoparticles are shown in fig 4.30 & fig 4.31. The peaks are characterized in table 4.13 & table 4.14 [56-65].

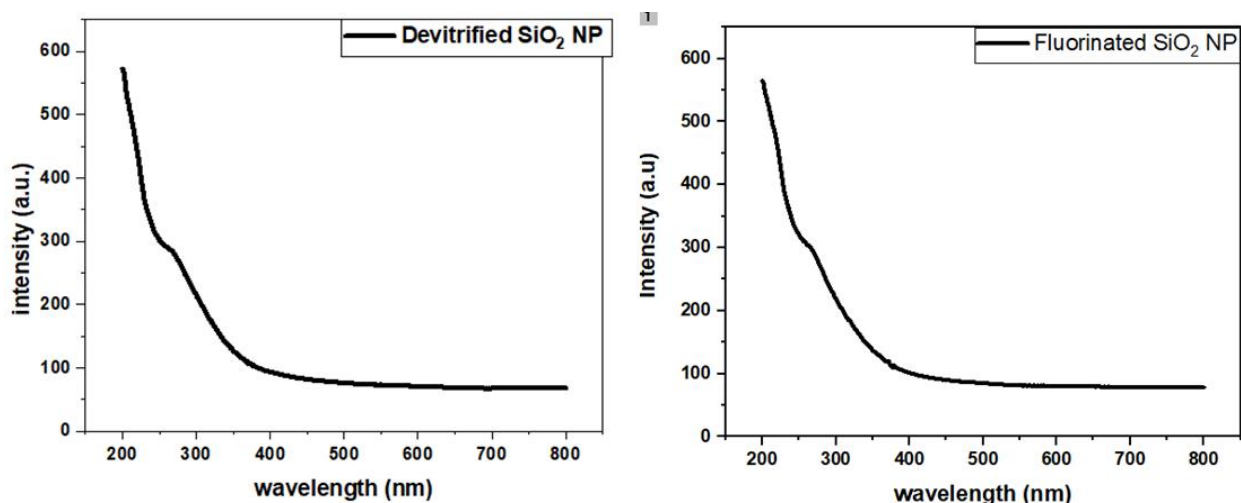


fig 4.28 UV Visible spectra of Devitrified & Fluorinated SiO₂ NPs.

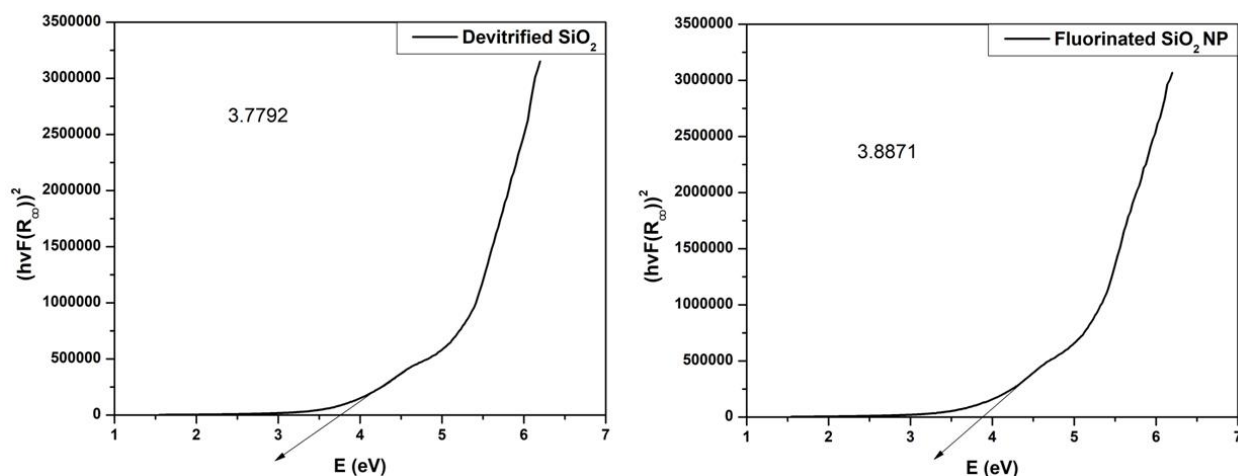


fig 4.29 Tauc plots of Devitrified & fluorinated nanoparticles respectively.

In the case of devitrified SiO₂ nanoparticles, the peaks of Si-O rocking vibrations, Si-O bending vibrations, Si-OH bending vibrations shifted to higher energies which may be because of heat

treatment. The absence of 952 cm^{-1} peak shows that Si-OH indeed has been vanished after devitrification due to dehydroxylation [47]. Also, the absence of free water peak (3414 cm^{-1}) shows that water has been vaporized due to high-temperature treatment. There is also the absence of C-H bonds.

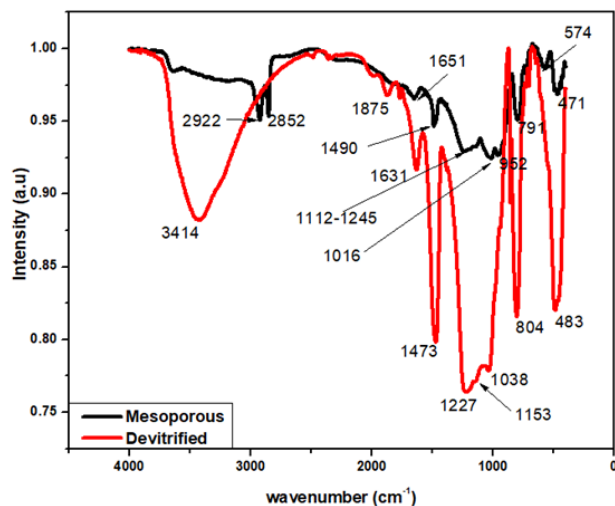


fig 4.30 FT-IR of Devitrified SiO_2 NP as compared to Mesoporous SiO_2 NP.

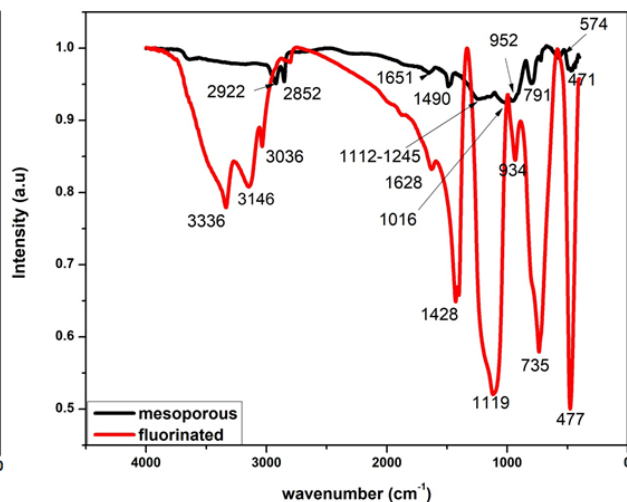


fig 4.31 FT-IR of Fluorinated SiO_2 NP as compared to original mesoporous one.

In the newly formed $(\text{NH}_4)_2\text{SiF}_6$ the new peaks have been originated shown in table 4.14 along with their positions.

M vs. H data as seen in fig 4.32 doesn't show any significant change. Now the devitrified sample can be doped with the anion to induce magnetism via pairing effect as stated earlier.

Bond nature	Peak position before devitrification (cm^{-1})	Peak position after devitrification (cm^{-1})
Si-O Rocking Vibrations	471	483
	574	
Si-O Bending Vibrations	791	804
Si-OH Bending Vibrations	952	
	1016	1038

Si-O-Si Asymmetric Vibrations	1112-1245	1227
	1490	1473
OH Bending	1651	1631
		1875
CH ₂ Symmetric Stretch [63-65]	2852	
-O-CH ₃ Stretching Vibrations	2922	
Free water peak	3414	

Table 4.13 IR peak positions before and after Devitrification.

Bond Nature	Peak position before Fluorination (cm ⁻¹)	Peak position after Fluorination (cm ⁻¹)
Free water Peak	3414	
Si-F stretching [56-57]		934
NH Stretching[58-60]		3336
Ammonia ions [61]		3146
Ammonium ions [61]		3036
NH ₄ ⁺ Bending mode [62]		1428

Table 4.14 Peak positions before and after fluorination.

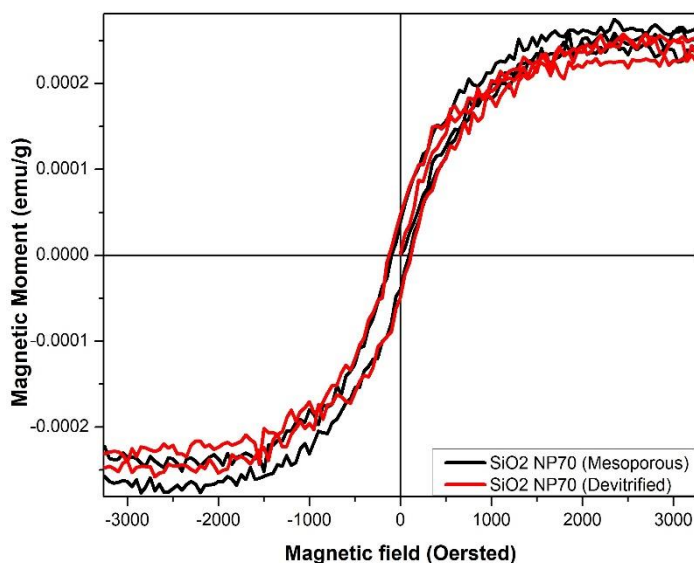


fig 4.32 M vs. H data of Mesoporous and Devitrified SiO₂ NPs.

4.10 Results & Discussions.

The devitrification hasn't affected the magnetic moment. But the future direction will be to dope with C or N or both to induce pairing effect between dopants and in turn to maximize the magnetic moment. The attempt of fluorination has failed because of the formation of ammonium fluorosilicate ((NH₄)₂ Si F₆).

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Chapter 5: Synthesis of doped SiO₂ nanoparticles.

5.1 Introduction.

As we know that vacancies and defects are somewhat correlated to magnetic properties of nanoparticles [1-8], we can induce these defects by artificially creating vacancies or introducing interstitials. The dopants will be Al, Mg, and Li. In the case of SiO₂ NP (440nm) doping with Al is done both on the surface and in bulk.

5.2 Synthesis of 2% bulk & surface doped SiO₂ NP (440nm).

We've used the same protocol as was followed for the synthesis of undoped nanoparticles [9] except making slight modifications to dope the nanoparticles [10-16]. The solution A was prepared by mixing 18.75 mL ethanol, 1.22 mL TEOS & 0.042 g of Al(NO₃)₃·9H₂O. The solution B was prepared by mixing 3.5 mL NH₄OH (25% in H₂O) 21.5 mL ethanol & 5 mL H₂O. The solution A was then added dropwise in Solution B with the rate of less than 0.05 mL/sec. The onset of turbidity was a clear indication of the formation of silica nanoparticles. The colloidal solution thus formed was allowed to age for 48 hours after which it was centrifuged and washed six times with ethanol after each runs and then dried at 60⁰C in a hot oven overnight.

The synthesis of 2% surface Al-doped nanoparticles was done with the same method as discussed earlier but with a difference that 0.042 g of Al(NO₃)₃·9H₂O was dissolved in sufficient amount of ethanol (0.1-0.5 ml) and was added after completely adding solution A in solution B [17-20]. The reason for adding it later in the reaction was to deposit most of the defect on the surface of nanoparticles contrary to the bulk case.

5.3 Synthesis of 5% Li, Mg & Al-doped SiO₂ NP (200nm).

The procedure employed was similar to the synthesis of SiO₂ NP (200nm) [21]. Required amount of TEOS & required amount of Al(NO₃)₃·9H₂O/Mg(NO₃)₂/LiNO₃ was added in 30 mL of EtOH and stirred vigorously for 5-10 minutes then it was allowed to sonicate for 10 minutes after which 1 mL distilled water was added, and the solution was again stirred for 5-10 minutes. Then

it was kept for sonication for 90 minutes, and then 2 mL of NH_4OH was added accompanied by stirring. The turbidity after 3-4 minutes was apparent and was a clear indication of the formation of nanoparticles. The turbid solution thus obtained was sonicated for 3 hours and aged overnight. After aging, it was centrifuged and washed several times with ethanol and dried for 1-2 days at 60°C in a hot air oven.

5.4 Characterization.

XRD, FT-IR, SEM, Diffuse reflectance UV & SQUID are the techniques employed in characterizations. The XRD of all the doped particles shows amorphous nature as can be seen in figs 5.1-5.2.

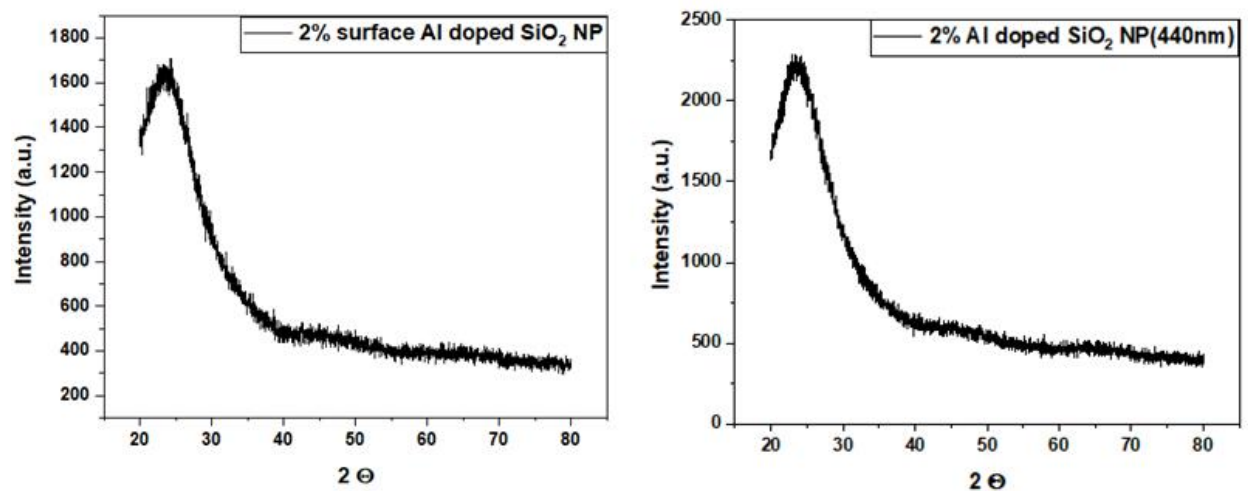
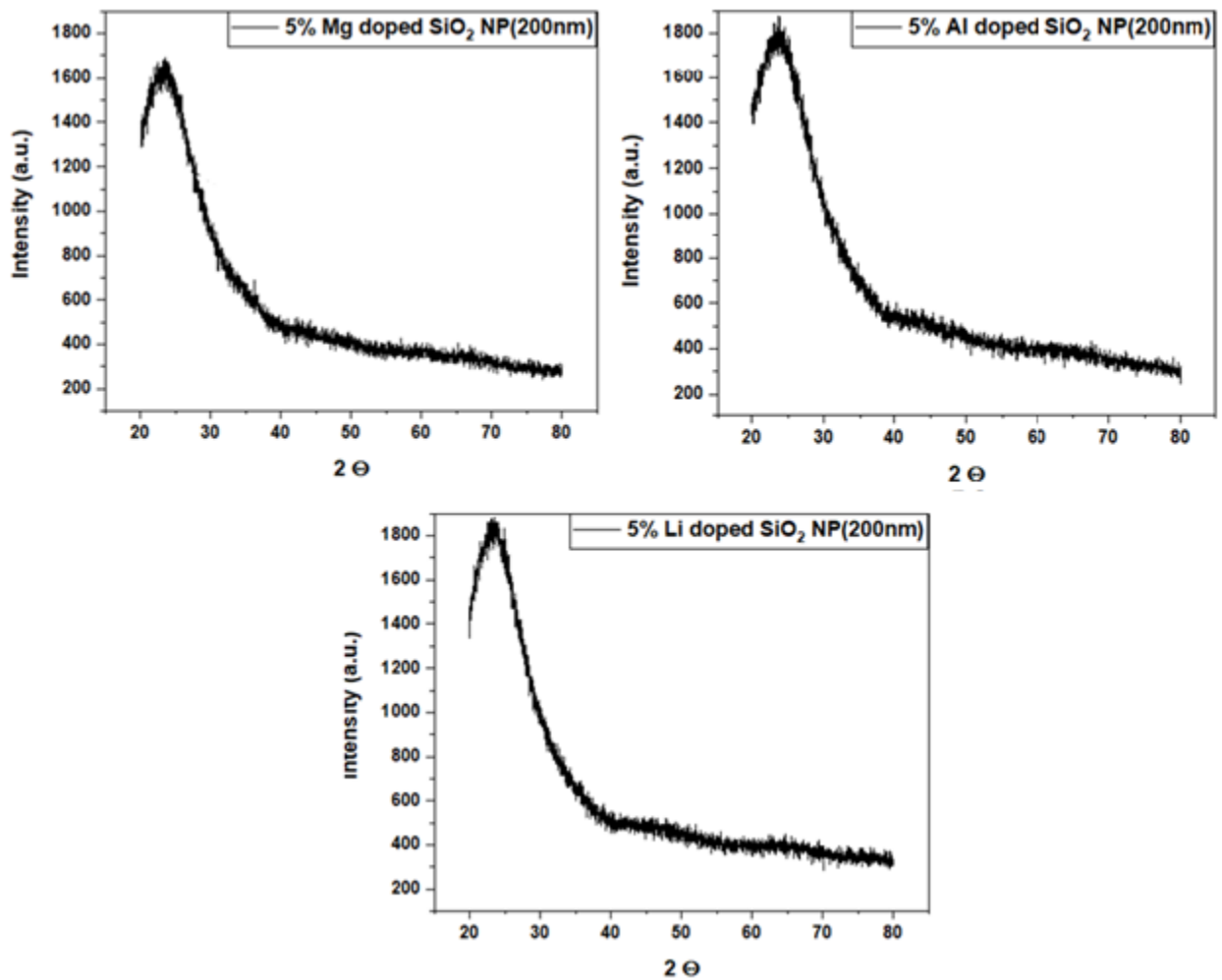


fig 5.1 XRD of 2% surface and bulk-doped (Al) SiO₂ NP (440nm).



Error! No text of specified style in document. **fig 5.2 XRD of 5% Mg, Al, Li-doped SiO₂ NP (200nm) respectively.**

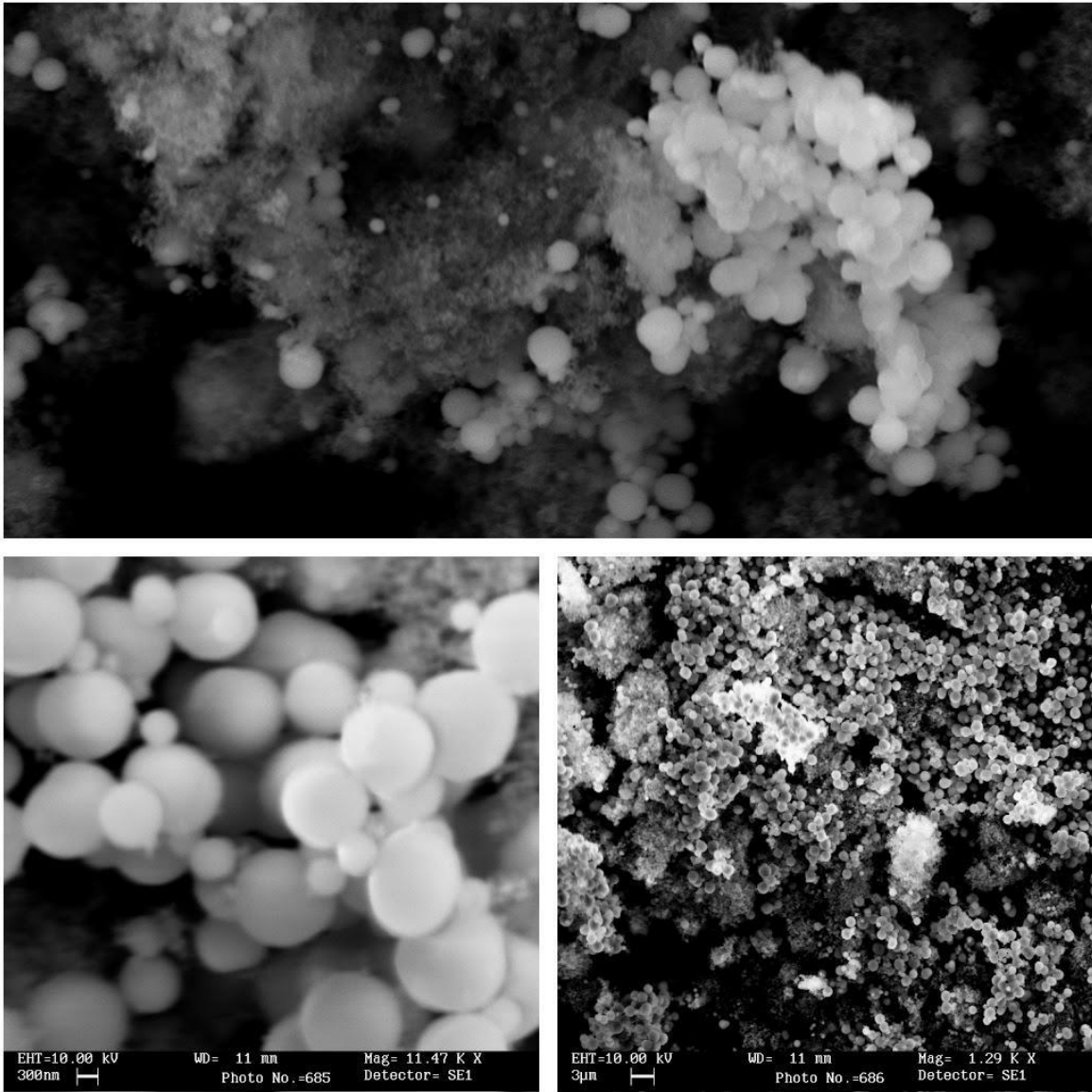


fig 5.3 SEM images of 2% Al-doped SiO₂ Nanoparticles.

As can be seen in figs 5.3-5.5 the nanoparticles are not monodispersed and uniform in size. The dopant salts which are added may be interfering with nuclei formation and growth. The SEM of both 5% Mg & 5% Li-doped SiO₂ NP (200nm) show non-spherical particles too in majority. The EDS (Energy dispersive x-ray spectroscopy) shows the presence of Mg peak (fig 5.7, Table 5.1). Even though elemental mapping shows presence of Mg, there is no clear way to tell if Mg cation is effectively doped into nanoparticles (fig 5.6).

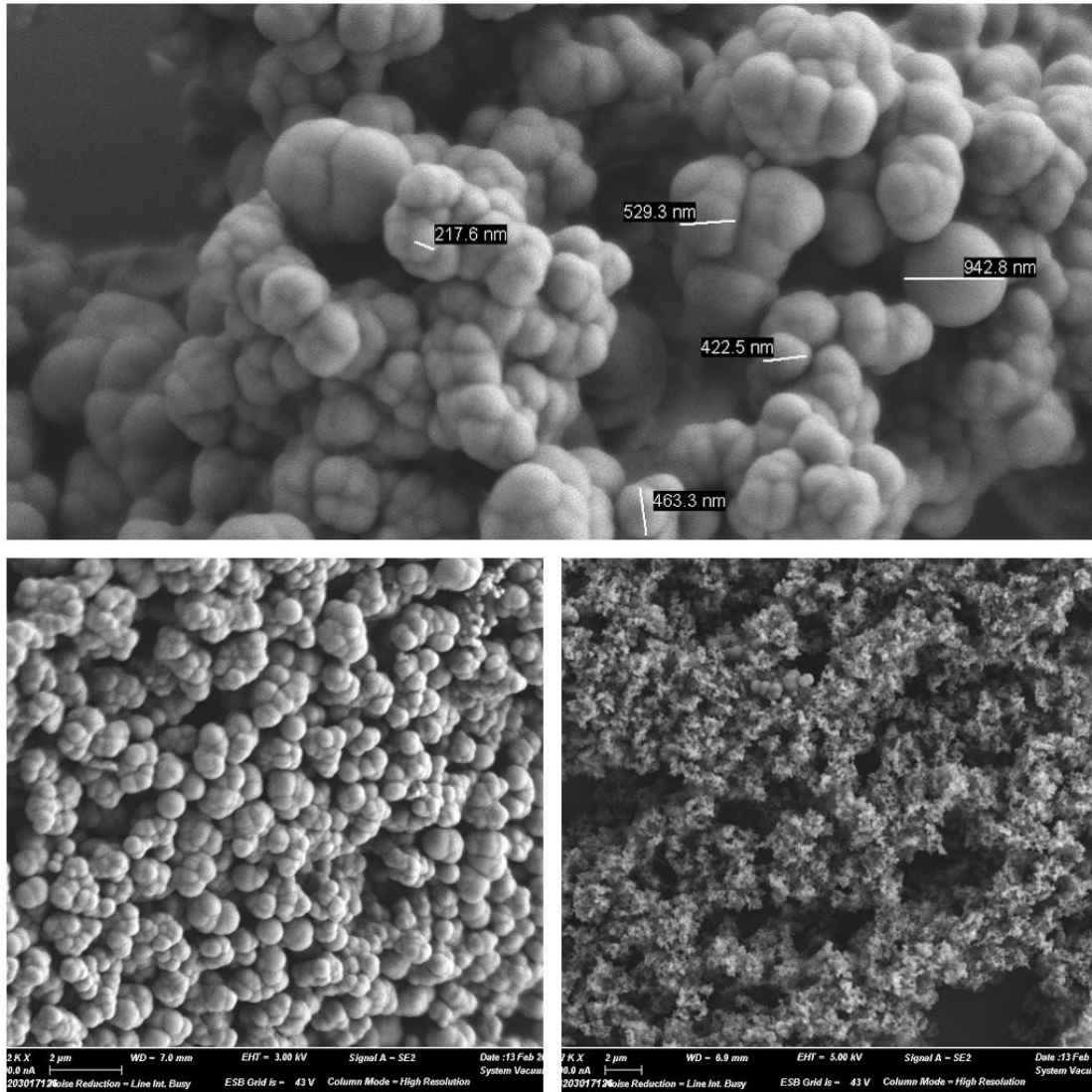


fig 5.4 SEM of 5% Li doped SiO₂ NP(200nm).

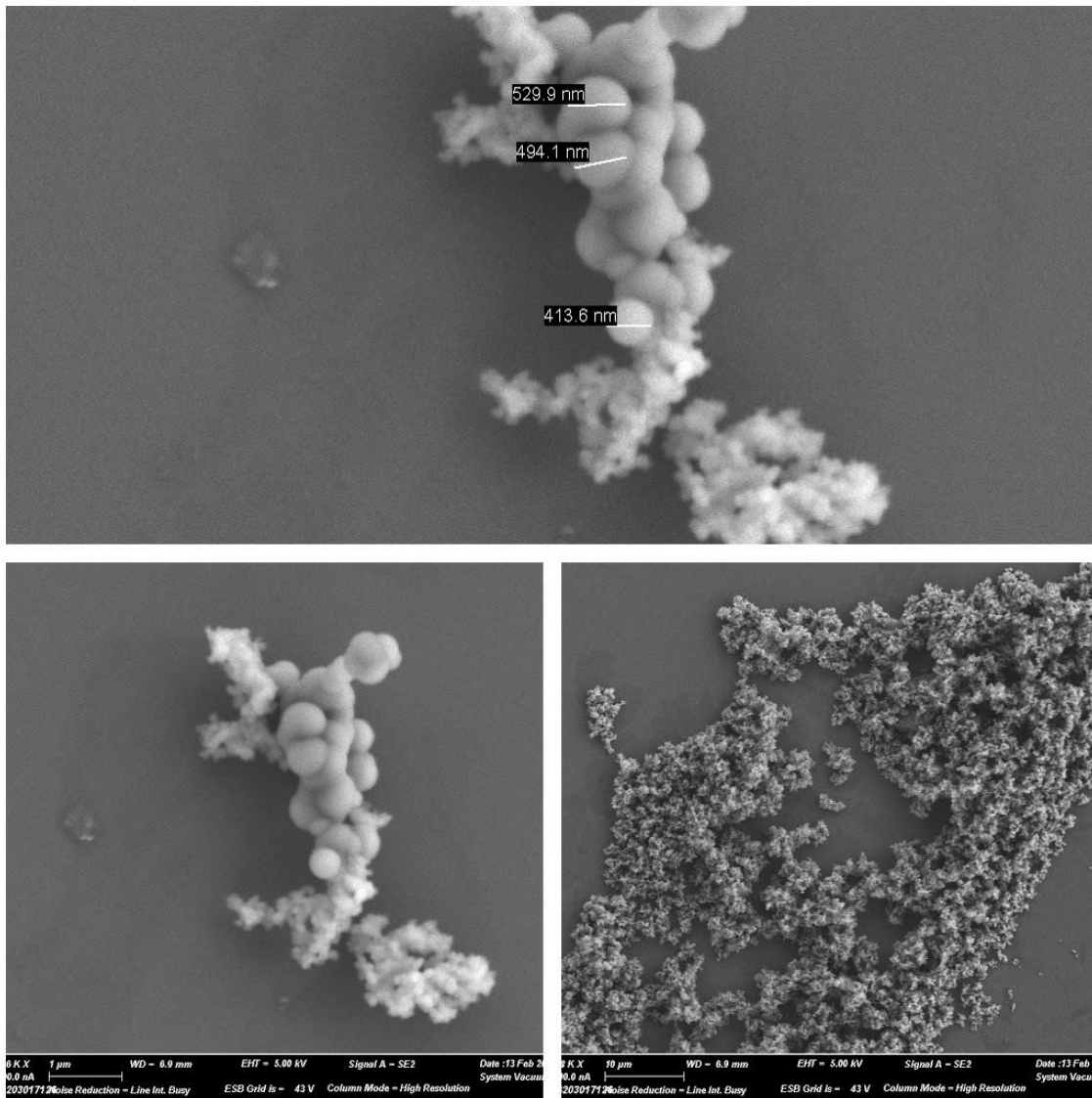


fig 5.5 SEM of 5% Mg-doped SiO₂ NP (200nm).

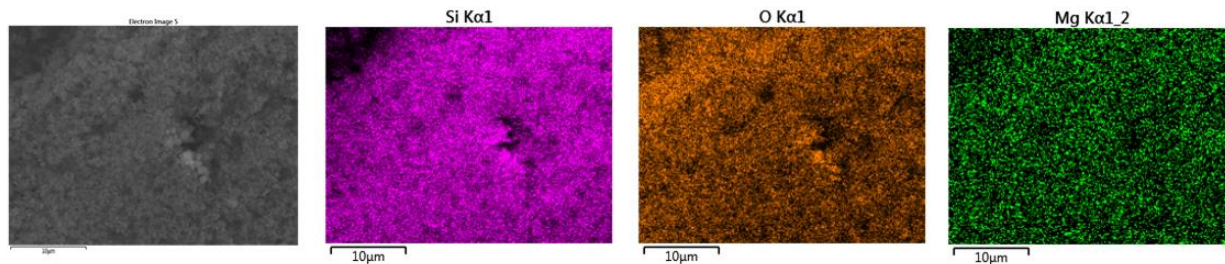


fig 5.6 Elemental Mapping of 5% Mg-doped SiO₂ NP (200nm).

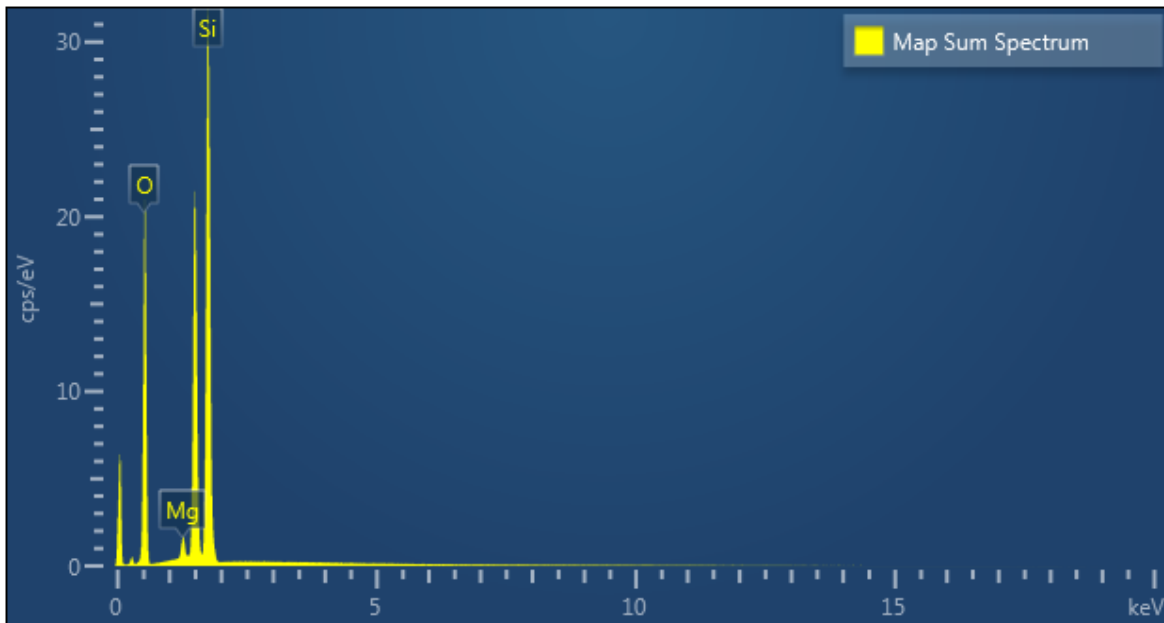


fig 5.7 EDS of 5% Mg-doped SiO₂ NP (200nm).

Element	Line Type	Apparent Concentration	k Ratio	Wt%	Wt% Sigma	Standard Label	Factory Standard	Standard Calibration Date
O	K series	36.46	0.12268	59.00	0.12	SiO ₂	Yes	
Mg	K series	1.08	0.00714	2.42	0.04	MgO	Yes	
Si	K series	20.07	0.15905	38.57	0.12	SiO ₂	Yes	
Total:				100.00				

Table 5.1 The result of EDS of 5% Mg doped SiO₂ NP (200nm).

M vs. H were measured to see the effect of doping on ferromagnetism. The higher magnetic moment of bulk Al-doped SiO₂ NP than the surface Al-doped SiO₂ NP may be because of more efficient doping or incorporation of the Al ions in the former case (fig 5.8). Since when we dissolve the salt of dopant in solution before the addition of TEOS; the dopant ions are already present in the solvent and can be incorporated in the nanoparticle in high numbers as compared to adding dopant salt after adding TEOS. M vs. H data (fig 5.9) of both Li & Mg-doped SiO₂ NP (200nm) shows six times less magnetic saturation than un-doped SiO₂ NP (200nm). This suggests that doping hasn't occurred efficiently.

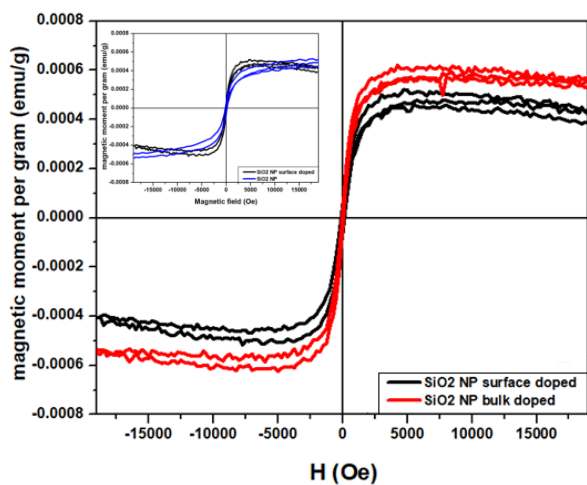


fig 5.8 M vs. H of surface & bulk Al-doped SiO₂ NPs (440nm). The box inside shows that surface Al doped & undoped particles have nearly the same magnetic moment.

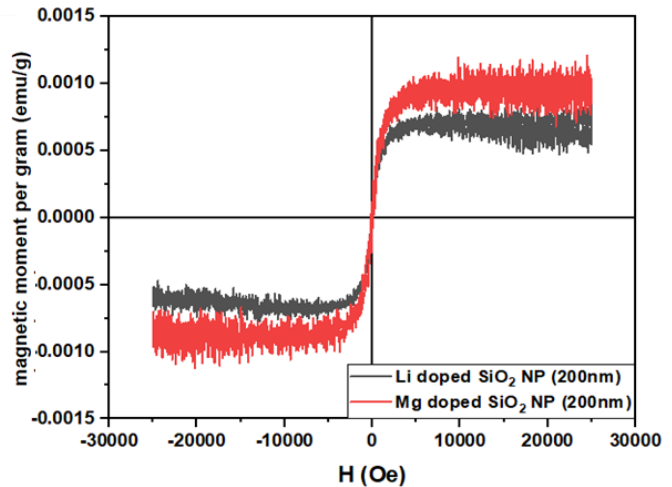


fig 5.9 M vs. H data of both Li & Mg-doped nanoparticles.

5.5 Results & Discussions.

Defects are induced more efficiently in the case of bulk Al-doped SiO₂ NP (440nm) than surface doped one. By looking at fig 5.9 & figs 5.4-5.6, we speculate that Li & Mg doping is too much inefficient meaning there is too less incorporation of dopant ions in nanoparticles. It seems that to enhance the magnetic properties of the nanoparticles; up to 10-15%, Al doping should be employed that too in sub 100nm particles to find their applications in the biomedical field [22].

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Chapter 6: Magnetism in BaTiO₃.

6.1 Introduction.

The materials that exhibit both magnetic & electric order are among the class of materials called multiferroic materials [1-10]. These types of materials are extensively used for data storage applications by assigning two opposing ferroelectric dipole moments or magnetic dipole moments as 0's and 1's [11]. They also find applications in magnetoelectric sensor, data storage, multifunctional materials, energy saving technologies & memory devices [12-17]. The multiferroicity till date has been found in boracites, fluorides, manganites of perovskites structures [18]. It is known that ferroelectric properties of BaTiO₃ nanoparticles diminish as the size reduces while oxygen vacancies & defects induced ferromagnetism increases [19]. But there are intermediate sizes at which these two properties can be successfully tuned making the BaTiO₃ nanocrystallites multiferroic [19]. The BaTiO₃ nanoparticles are thus synthesized. The samples are also annealed at several temperatures. The cubic BaTiO₃ nanoparticles (non-ferroelectric) were also synthesized. SEM, XRD, EDS, SQUID characterizations were also done.

6.2 Synthesis of cubic BaTiO₃ nanoparticles.

The cubic BaTiO₃ was prepared by precipitation reaction [20]. 1g of NaOH was added to 5 mL of 0.2 mmol/L Ba(NO₃)₂ solution with continuous stirring and heating in an oil bath. After the temperature of the oil bath reached 80⁰C, 0.3 mL titanium isopropoxide was added, and the resultant white colored turbid solution was kept for stirring for 24 hours after which it was centrifuged and washed with distilled water 3 times after each run. The as-obtained powder was later dried in a hot air oven.

6.3 Synthesis of tetragonal BaTiO₃ nanoparticles.

Tetragonal BaTiO₃ nanoparticles were synthesized by the polymeric precursor process [19]. 180 mL ethanol & 60 mL acetic acid was mixed to which 0.88 mL of titanium isopropoxide (Ti(OR)₄) was added and kept for stirring for 2 hours to form solution A. To prepare solution B 0.592 g Ba(CO)₃ & 2.3054 g citric acid were dissolved in 20 mL de-ionized water. Then after completion of stirring solution A, solution B was added and immediately after the addition of

solution B poly vinyl alcohol (PVA) solution (2.5 g PVA dissolved in 96.35 mL) was added. The onset of turbidity was an indication of the formation of nanoparticles. The solution was then centrifuged and was dried in a hot air oven at 60⁰C for more than 1-2 days. And the dried powder was calcined in air at 700⁰C for 1 hour to remove the organic materials.

6.4 Characterization.

The XRD peaks are without any impurity as evident in fig 6.1. tetragonal BaTiO₃ was annealed at several temperatures. The annealing temperature is indicated as BTO700 (for 700⁰C) & so on for different temperatures. The EDS as well as elemental mapping of both BTO 1000 & BTO 1300 confirms the presence of Ba, Ti & O as shown in figs 6.4-6.5 & figs 6.7-6.8 & tables 6.1 & 6.2. The SEM of BTO1300 shows increased in particles size (than BTO1000) which is due to the annealing.

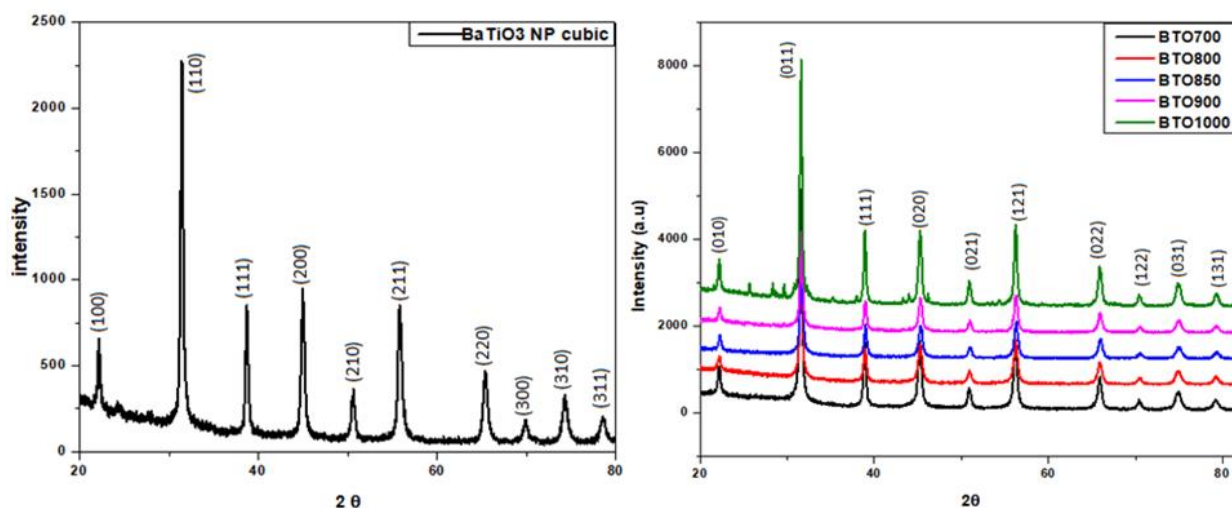


fig 6.1 XRD of Cubic & Tetragonal BaTiO₃. Latter was annealed at several temperatures.

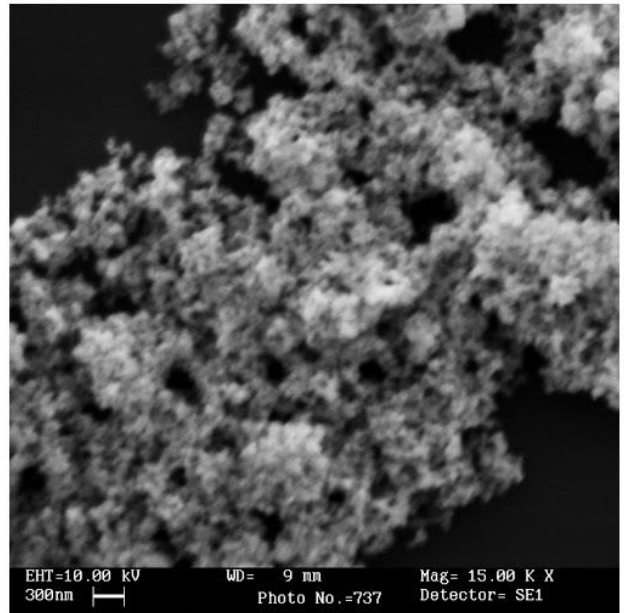
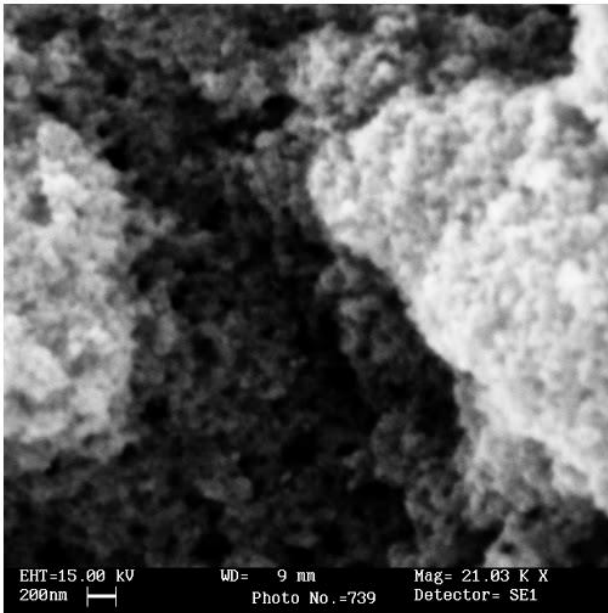
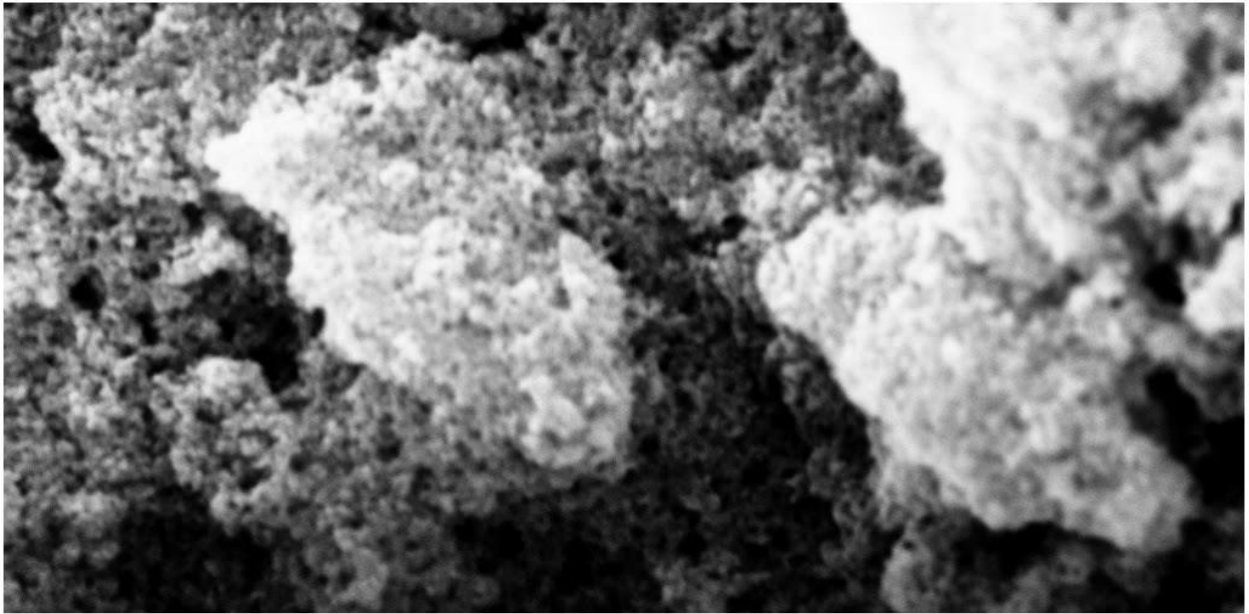


fig 6.2 SEM of BTO 700.

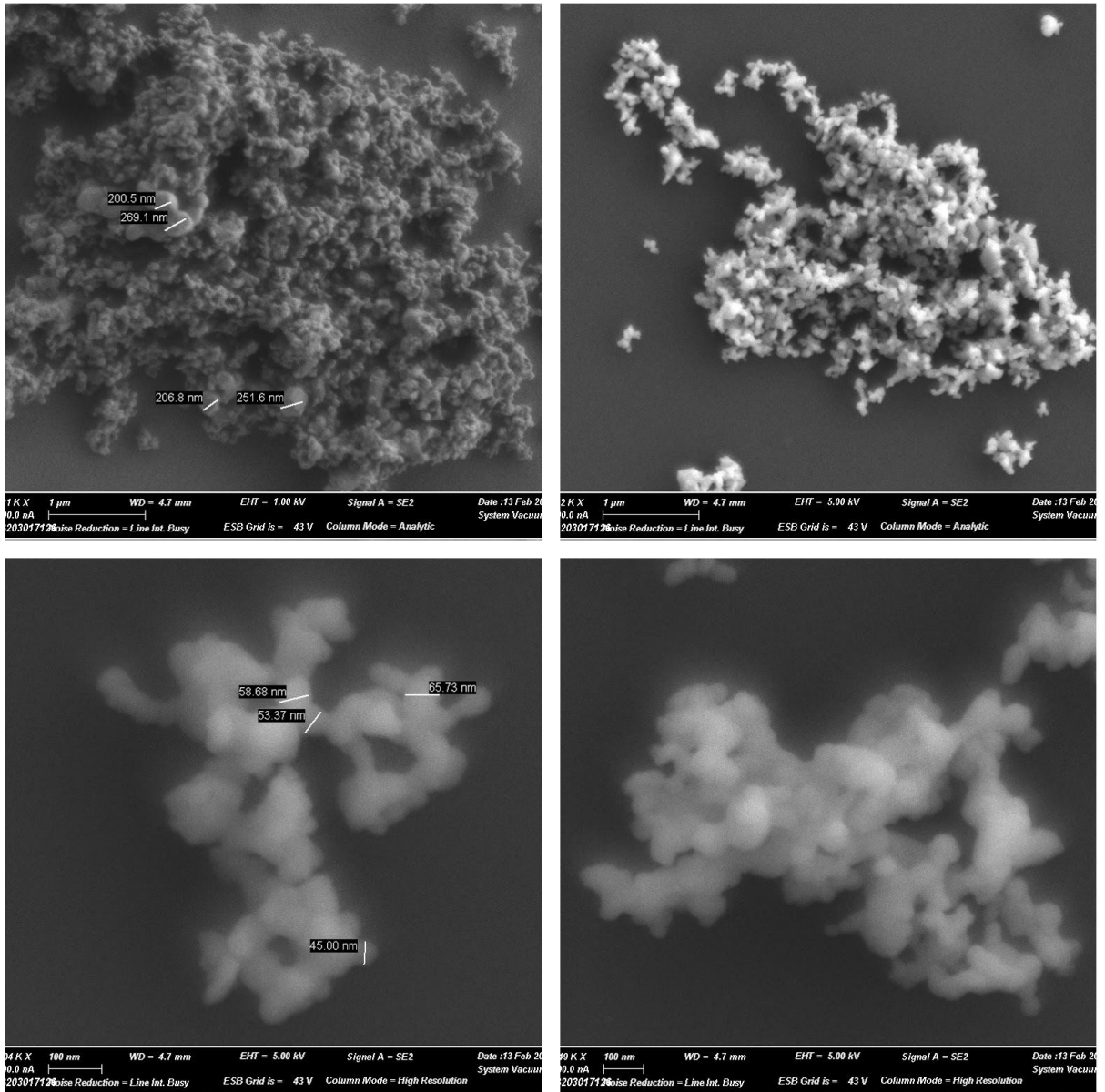


fig 6.3 SEM images of BTO 1000.

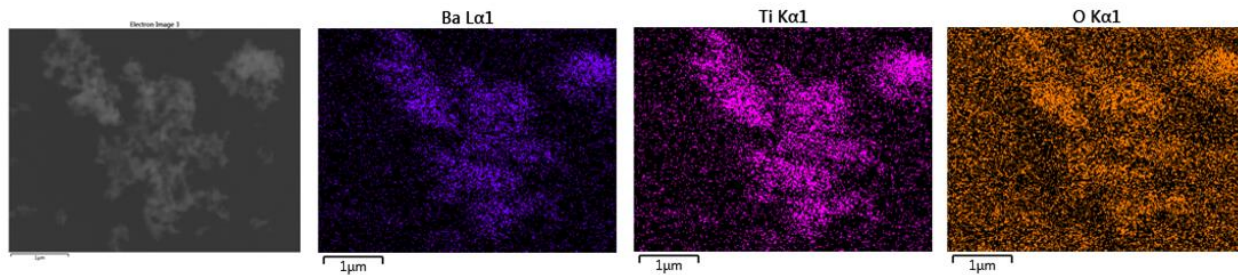


fig 6.4 Elemental mapping of BTO 1000.

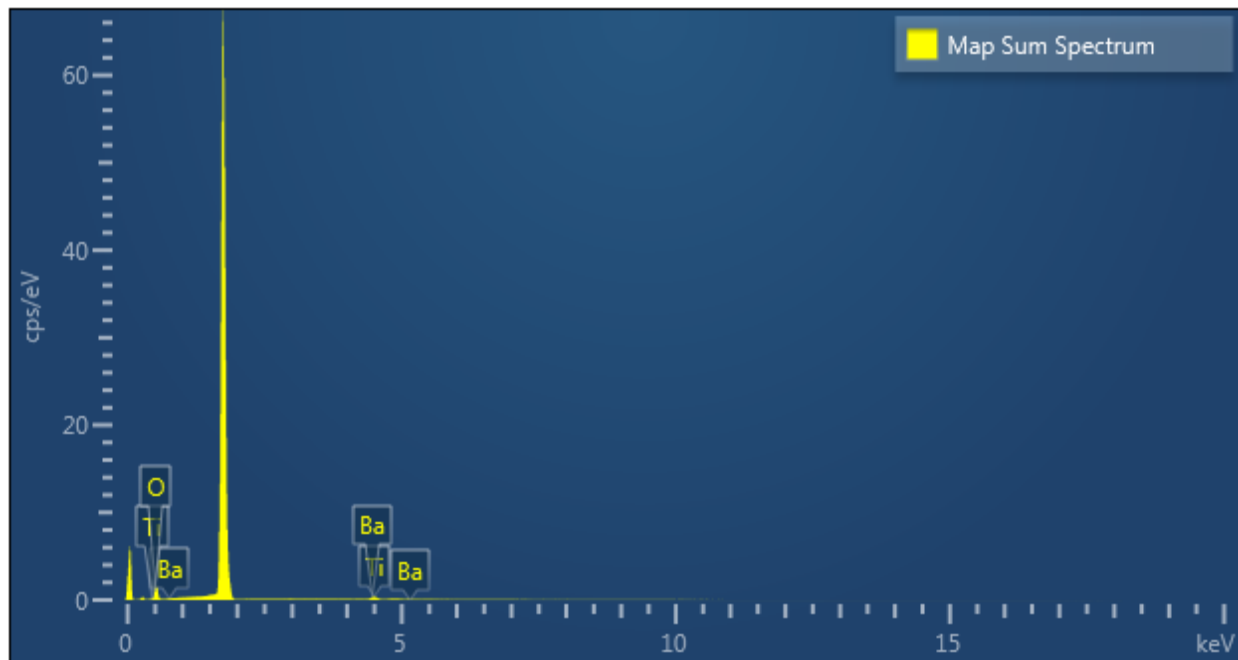


fig 6.5 EDS of BTO 1000.

Element	Line Type	Apparent Concentration	k Ratio	Wt%	Wt% Sigma	Standard Label	Factory Standard	Standard Calibration Date
O	K series	2.93	0.00987	70.27	0.82	SiO2	Yes	
Ti	K series	0.27	0.00272	9.79	0.42	Ti	Yes	
Ba	L series	0.53	0.00496	19.94	0.83	BaF2	Yes	
Total:				100.00				

Table 6.1 Elemental composition (by wt%) of BTO 1000.

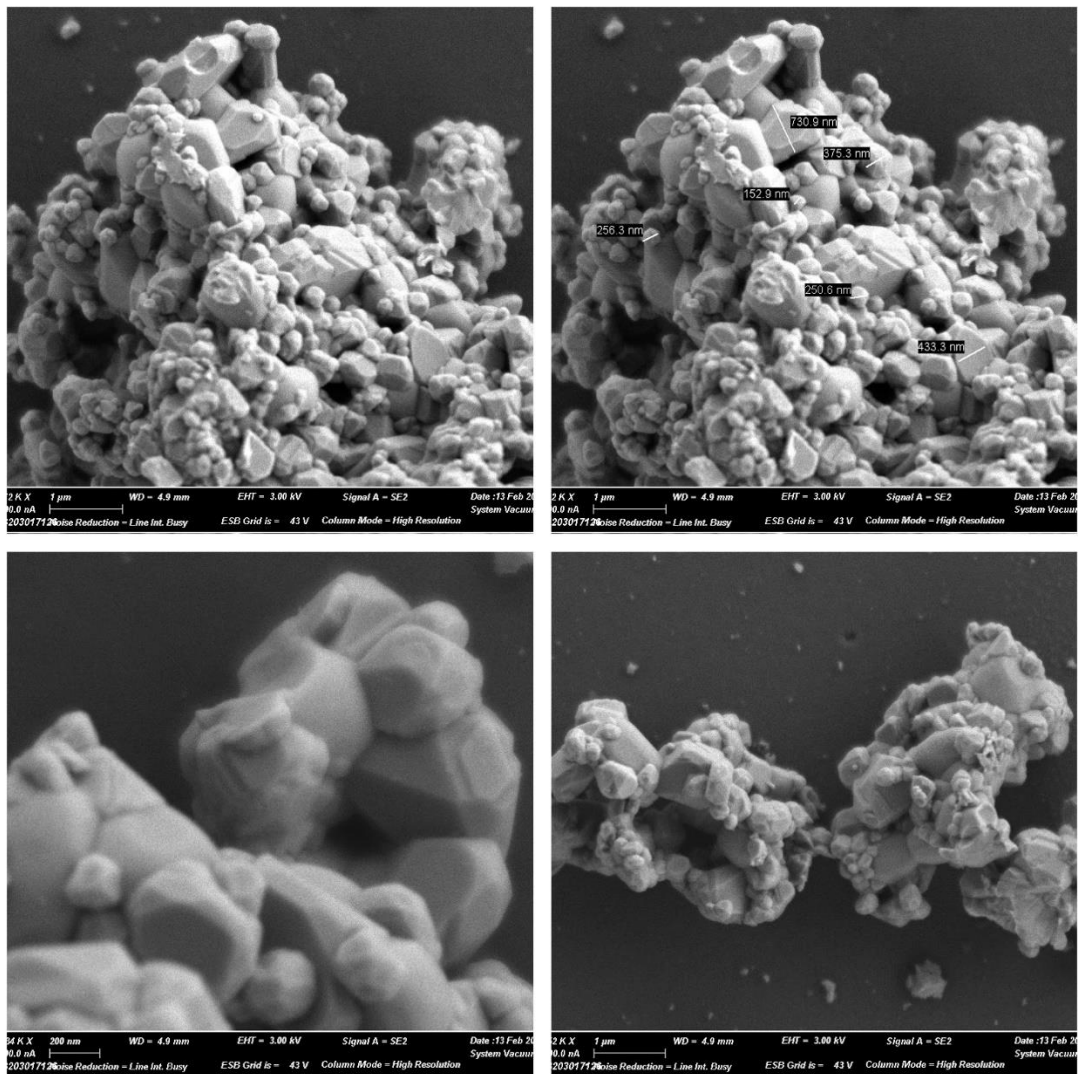


fig 6.6 SEM of BTO 1300.

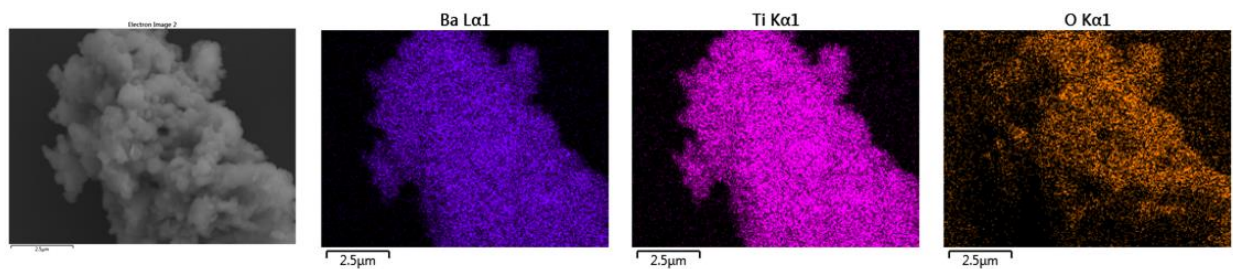


fig 6.7 Elemental Mapping of BTO 1300.

Element	Line Type	Apparent Concentration	k Ratio	Wt%	Wt% Sigma	Standard Label	Factory Standard	Standard Calibration Date
O	K series	5.36	0.01803	28.54	0.33	SiO2	Yes	
Ti	K series	5.00	0.05005	21.62	0.23	Ti	Yes	
Ba	L series	11.05	0.10348	49.83	0.33	BaF2	Yes	
Total:				100.00				

Table 6.2 Elemental composition (by wt%) of BTO 1300.

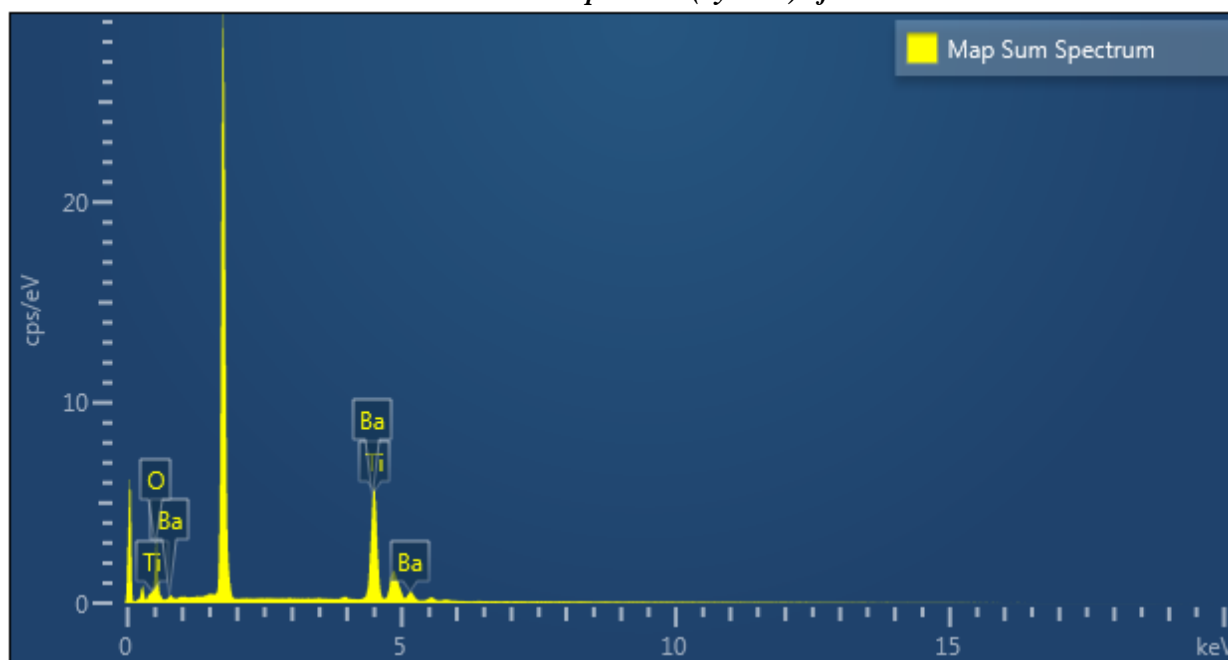


fig 6.8 EDS of BTO 1300.

As can be seen in fig 6.9 the tetragonal BaTiO₃ shows much more saturation magnetic moment than the cubic BaTiO₃. The saturation of both cubic & tetragonal nanoparticles are 0.0004 emu/g & 0.007 emu/g respectively. Annealing has reduced the magnetic moment by seven times.

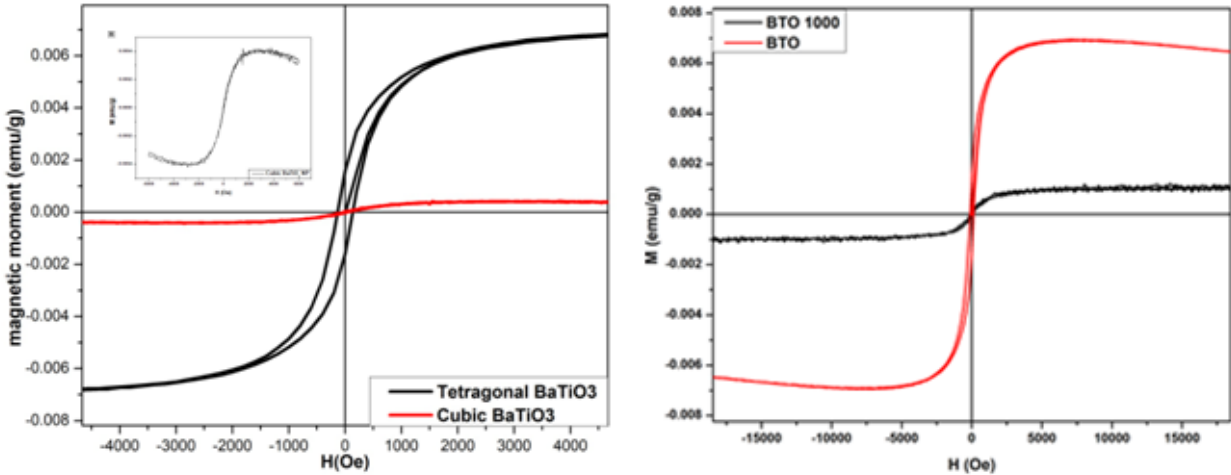


fig 6.9 M vs. H data of both Tetragonal & Cubic BaTiO₃ (left). The inside box shows M vs. H of cubic BaTiO₃. While on right M vs. H data before & after annealing is shown.

SANS (small angle neutron scattering) data of BaTiO₃ nanoparticles is shown in fig 6.10. The data taken is $I(Q)$ vs. Q . I is intensity & Q is the wave vector. If the nanoparticles were simple the, $I(Q)$ vs. Q would have been a straight line (Porod's law) [21], but it's not the case. The line is humped. The hump is due to magnetic scattering. Magnetic scattering is the scattering due to the interaction of the magnetic moment of neutrons with the electronic moments in the sample [22]. The humped line is fitted to a spherical shell model which indicates that the core is nonmagnetic and the shell is magnetic. Thus the magnetism indeed exists on the nanoparticle surface. This is in agreement with the previous studies [23-29].

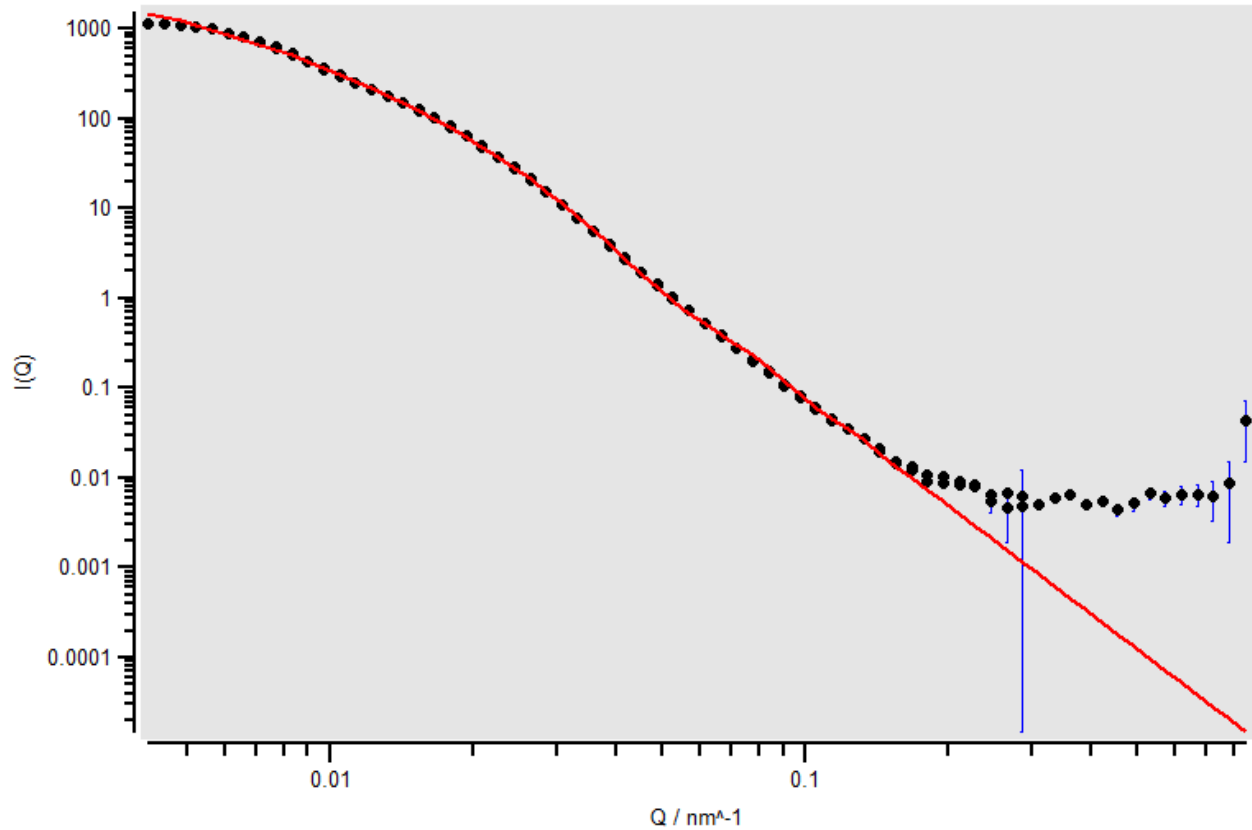


fig 6.10 I vs. Q data of BaTiO₃ Nanoparticles fitted as a Spherical Core Shell Model.

The particle size distribution is a lognormal distribution (fig 6.12) with an average core of 44.2 nm & the average shell of 10 nm. The shell's 10 nm size corresponds to magnetic order along with its actual physical size.



fig 6.11 The structure of BaTiO₃ can be visualized as dark nonmagnetic core (99nm) and outer shell exhibiting magnetic order (fuzzy surface).

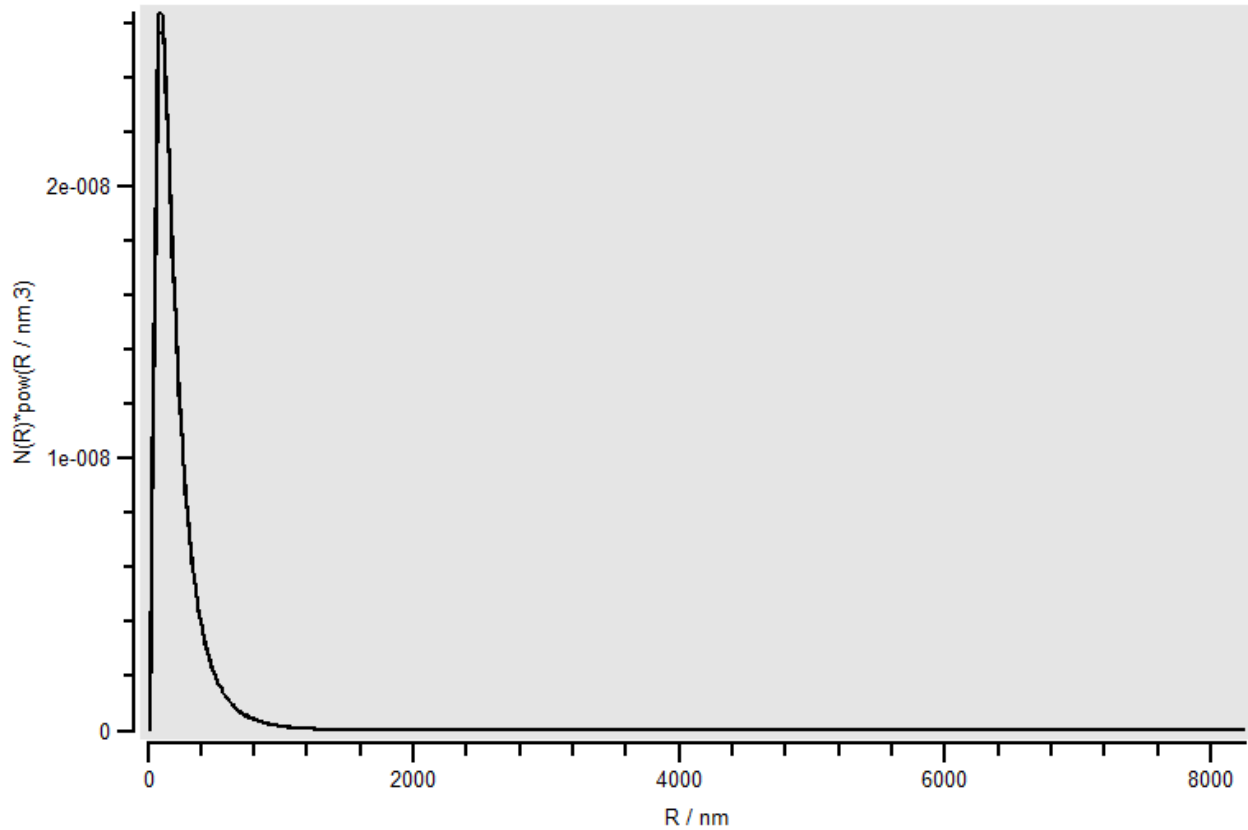


fig 6.12 Lognormal distribution of the core size of nanoparticles.

6.5 Results & Discussions.

The as-synthesized Tetragonal BaTiO₃ nanoparticles show nearly seven times more saturation than that of previously reported [19] although the same synthesis route was followed. The reason for this may be due to annealing carried out in the air rather than oxygen (the latter may have been detrimental to magnetic moment). The tetragonal BaTiO₃'s higher magnetic moment than the cubic BaTiO₃ may be due to its more specific surface area than cubic BaTiO₃. The tetragonal BaTiO₃ is stable until 1000⁰ C of heating. The ferromagnetism (in tetragonal BaTiO₃) has originated from the surface as can be seen from the analysis of SANS data. The annealing in the air brings down the magnetic moment which may be due to a decrease in concentration of oxygen vacancies.

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Chapter 7: Conclusion & Outlook.

7.1 Conclusion.

The study to understand the ‘origin of ferromagnetism in otherwise nonmagnetic oxides’ was focused on two systems i.e. nanoparticles of SiO_2 & BaTiO_3 . In both the cases we are successful in inducing a reasonable amount of magnetic moment i.e. as high as 0.006 emu/g. The SiO_2 system has proven contradictory in case of the mesoporous SiO_2 nanoparticles which show more oxygen vacancies (as seen in photoluminescence spectra) but show lesser magnetic moment than SiO_2 nanoparticles of 440nm size (which exhibits too less oxygen vacancies than mesoporous SiO_2 nanoparticles). SiO_2 nanoparticles with size 200nm show about six times more magnetic moment than those with size of 100nm. The higher magnetic moment in 200nm sized SiO_2 nanoparticles is because of their different mode of syntheses which involved ultrasonication of about 3 hours and 90 minutes. In synthesis of SiO_2 NP (200nm); after the initiation of nuclei formation (i.e. appearance of turbidity of solution) the solution was kept for 3 hours under continuous sonication which has resulted in more oxygen vacancies & defects on the surface of nanoparticles and thus increased the magnetic moment. This result is in agreement with the studies of Jorge Mejia et al. which has compared effect of ultrasonication on nanoparticle surface as opposed to stirring [4].

Doping of Al in SiO_2 nanoparticles proves to be very efficient way to enhance ferromagnetism than that of Li or Mg doping. The latter two cations were harder to incorporate in SiO_2 nanoparticles in first place.

Annealing the nanoparticles decreases oxygen vacancies and thus magnetic moment in both SiO_2 & BaTiO_3 cases.

The band gap of SiO_2 nanoparticles decreased after annealing and also with decreasing size.

Small angle neutron scattering (SANS) analysis of BaTiO_3 nanoparticles shows that the core is nonmagnetic while the surface has magnetic moments which confirms that the origin of ferromagnetism is indeed due to surface defects.

7.2 Future Work.

The current SANS data of BaTiO₃ (taken under incidence of unpolarized neutron beam and in absence of magnetic field) shows only the evidence of presence of magnetic structure. But the same measurements should be done in presence of magnetic field and using polarized neutron source in order to obtain the magnetic structure [1]. Similarly for SiO₂ NP (200nm) the SANS data should be taken in order to analyze its magnetic structure.

The ion beams can be made incident on SiO₂ nanoparticles to induce defects and vacancies and thus enhancing the ferromagnetism.

Since we have succeeded in devitrification of SiO₂ nanoparticles; the anion doping (e.g. C or N) can be tried in order to increase the magnetic moment of SiO₂ nanoparticles (i.e. by inducing pairing interactions between dopants) [2]. Fluorination is another way by which ferromagnetism can be enhanced.

Since SiO₂ NP (200nm) (which was prepared by employing ultrasonication) shows more magnetic moment than SiO₂ NP (100nm) the effect of ultrasonication on the surface of nanoparticles should be studied in more detail using XPS (x-ray photoelectron spectroscopy), TEM (transmission electron microscopy) etc. If the ultrasonication is utilized for the synthesis of sub 100nm nanoparticles the magnetic moment will be highly enhanced.

Muon spin resonance spectroscopy (μ SR) which can probe local magnetic moment can be used to investigate this type of ferromagnetism.

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